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Multiple Scattering Corrections for the Associated-Particle Neutron Time-of-Flight Technique

**U.S. DEPARTMENT OF COMMERCE
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Abstract

The computer code, MAGGIE, for the calculation of multiple scattering and sample attenuation in neutron differential cross-section measurements, has been revised and corrected. The particular case of the scattering geometry required by the associated-particle time-of-flight is considered in detail.

Key words: associated-particle, cross-section,
Monte Carlo, multiple scattering,
neutron, time-of-flight

MULTIPLE SCATTERING CORRECTIONS FOR THE ASSOCIATED-PARTICLE NEUTRON TIME-OF-FLIGHT TECHNIQUE

Allan C. B. Richardson

I. INTRODUCTION

Measurements of fast neutron elastic and inelastic differential cross sections have, for many years now, usually been done using one of two time-of-flight techniques. The first of these requires a pulsed source of neutrons, and energy separation of the various neutron groups is then achieved by time correlation of the scattered neutrons with the incident neutron pulse. This technique has the advantage of flexibility of incident neutron energy and intensity, and the disadvantages of a relatively high time-correlated background and a low duty cycle. The other, and less commonly used, technique utilizes the detection of a charged particle from the neutron source reaction to tag the incident neutrons in time and direction. Energy separation of the various scattered neutron groups is then achieved by time correlation with the incident neutrons. This "associated-particle technique" has the advantages of very small time-correlated background, inherent absolute determination of the incident correlated neutron flux, and high duty cycle; but it suffers from limitations on the available neutron intensity and energy. However, at energies and intensities where this technique is applicable, it is the method of choice, since it is capable of yielding results of high accuracy without the ambiguities introduced by the time-correlated backgrounds and the massive shielding required by pulsed source techniques. The source reactions eligible for this method are those involving very light nuclei, and therefore capable of producing a light (and thus energetic) stable recoil nucleus. The $T(d,n)^4\text{He}$ reaction, producing 14-15 MeV neutrons at 90° over a rather wide range of incident particle energies, is most commonly used; other reactions that have been employed are $D(d,n)^3\text{He}$ and $T(p,n)^3\text{He}$, both of which produce lower energy neutrons. We will confine the discussion here to the $T(d,n)^4\text{He}$ case, although the method described is more generally applicable.

The associated-particle technique has scattering sample requirements that are quite different from those for a pulsed source. Instead of a relatively uniform incident neutron flux across the sample, the correlated neutron beam is highly directional. The angular distribution about the neutron beam line is usually well approximated by

$$I = I_0 e^{-\left(\frac{\theta}{\theta_0}\right)^2},$$

with θ_0 typically only a few degrees [1]. This property can be very useful [2]. It provides a high degree of neutron collimation without the need for massive collimators, which, in the case of 14 MeV neutrons, can produce substantial degradation of the initially monoenergetic neutron

beam. However, it forces the use of scattering samples of uniform thickness so that the cylindrical or spherical samples commonly used to simplify multiple scattering corrections are immediately ruled out. Otherwise a detailed knowledge of the neutron beam shape and extremely accurate alignment of this beam with respect to the scattering sample is needed. This alignment problem is further complicated because the center of the neutron beam slowly moves back as the neutron producing target ages with use. In measuring angular distributions of scattered neutrons it is of course also desirable to reduce the amount of scattering sample not directly in the neutron beam, so as to minimize multiple scattering.

A scattering sample in the form of a truncated cone, axis lying along the neutron beam, best satisfies all of these requirements. The origin of this cone is taken sufficiently far behind the source to allow for finite spot size on the neutron producing target, to provide some flexibility in alignment, and also to make allowance for changes due to target aging during a run. In order to make best use of the available neutron intensity, one must also use the thickest sample possible. The limit is set by either the angular resolution required at 90° or the time resolution required. A typical geometry is shown in figure 1.

In either case the resulting samples are sufficiently thick to require a careful multiple scattering correction. None of the analytical techniques [3], useful at energies up to a few MeV, are adequate at 14 MeV, the energy most commonly used for measurements of this type, since at energies above 6-7 MeV the diffraction peaks in the elastic angular distributions become too numerous. The only method of sufficient generality is the Monte Carlo technique. A survey of existing Monte Carlo codes revealed none for this particular geometry, but it was immediately apparent that the code "MAGGIE," developed by Parker, Towle, Sams, et al. at Aldermaston [4,5] contained all of the other elements important to such a calculation. For example, this code easily accommodates the neutron source distribution specified above. In addition, MAGGIE calculates an energy spectrum at each angle, so that false peaks due to double scattering can easily be identified. This sophistication is often useful, for example: the measurement of inelastic scattering from the very weakly excited 7.66 MeV 0^+ level in ^{12}C is easily confused by double scattering from the more easily excited 4.43 MeV 2^+ level in this nucleus.

A corollary need of any Monte Carlo neutronic calculation is an easily accessible, but at the same time sufficiently general, file program for the nuclear data required. This was available in a companion code to MAGGIE, entitled MOULD [6]. We have therefore modified the code MAGGIE so that it now handles the geometry required for associated-particle time-of-flight measurements. During the course of this modification we also corrected a few coding errors found in the original version, and made several additional modifications, principally updating the code to current computer syntax and capabilities.

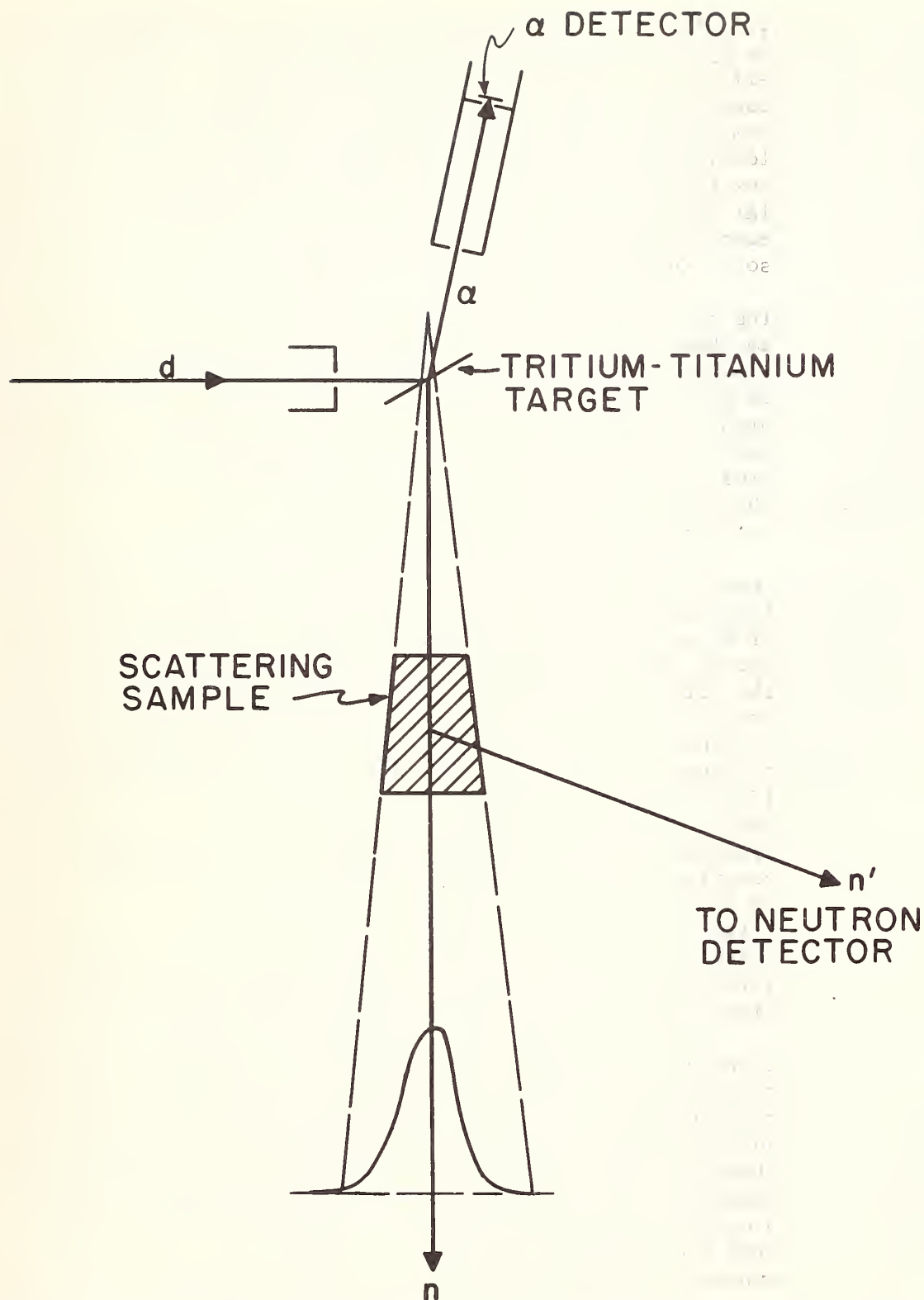


Fig. 1. The experimental geometry for associated-particle measurements at 14 MeV. The associated neutron intensity profile is indicated at the bottom of the figure.

Although a complete description of the code would be out of place here (see references 4 and 5 for details), a general overview of the program will facilitate later discussion of specific details. The point of view adopted is first to sample the scattering sample geometry and the available nuclear interactions using straightforward Monte Carlo techniques. Then, at each collision point, the weighted probability of scattering and escape to each of 33 detector positions (angles) in a half-plane lying to one side of the neutron beam axis is computed using the experimental data for elastic events and each of those inelastic processes of interest. The use of weighted probabilities of scattering to each of the detector positions at the final collision in the sample in place of a completely Monte Carlo approach results in a greatly reduced computation time. This is because the small solid angle subtended by each of the detectors (typically 10^{-3}) makes a final Monte Carlo scattering particularly inefficient. The output angular distributions obtained from applying this procedure to, typically, 1000 interacting neutrons are then reflected about the experimental input data and used as input for a second iteration. Two or three iterations are usually sufficient to obtain a convergent result.

In section II we give a description of the changes made in the code. First, the new coding for the scattering geometry used in the associated particle technique is described. Next we discuss a number of coding errors in the original version. Finally some changes are described that simplify the code and adopt it to FORTRAN V syntax. In section III we give some results obtained using data from the scattering of 14 MeV neutrons on natural carbon. A listing of those subroutines entirely rewritten or having extensive changes is given in Appendix A. In figure 2 the calling sequence for all of the components of the program is shown. A brief description of all of these subroutines appears in Appendix B.

II. DESCRIPTION OF MODIFICATIONS TO THE MONTE CARLO CODE MAGGIE

A. Changes Due to the New Scattering Sample Shape

Four subroutines are affected by changing the shape of the scattering sample. These are: 1) subroutine INPUT - those sections where scatterer parameters are read in and the flux attenuation factor is calculated are affected. The flux attenuation factor is defined as the ratio of incident flux along the axis of the sample to average flux in the entire sample. 2) CRNEU, the subroutine that creates incident neutrons at the entrance face of the scatterer by random sampling of the incident neutron spatial distribution. 3) TRACK, the tracking subroutine. 4) FPATH, the subroutine for calculating the probability of neutron escape from the sample in the direction of each of the assumed detector positions, for each collision point arrived at in TRACK.

It is worth noting here, although not necessary to what follows, that the data used for a) the neutron track lengths in the Monte Carlo sampling of the sample shape from subroutine EGMV, and b) the Monte Carlo sampling of reaction type and angular distribution at each collision by subroutine

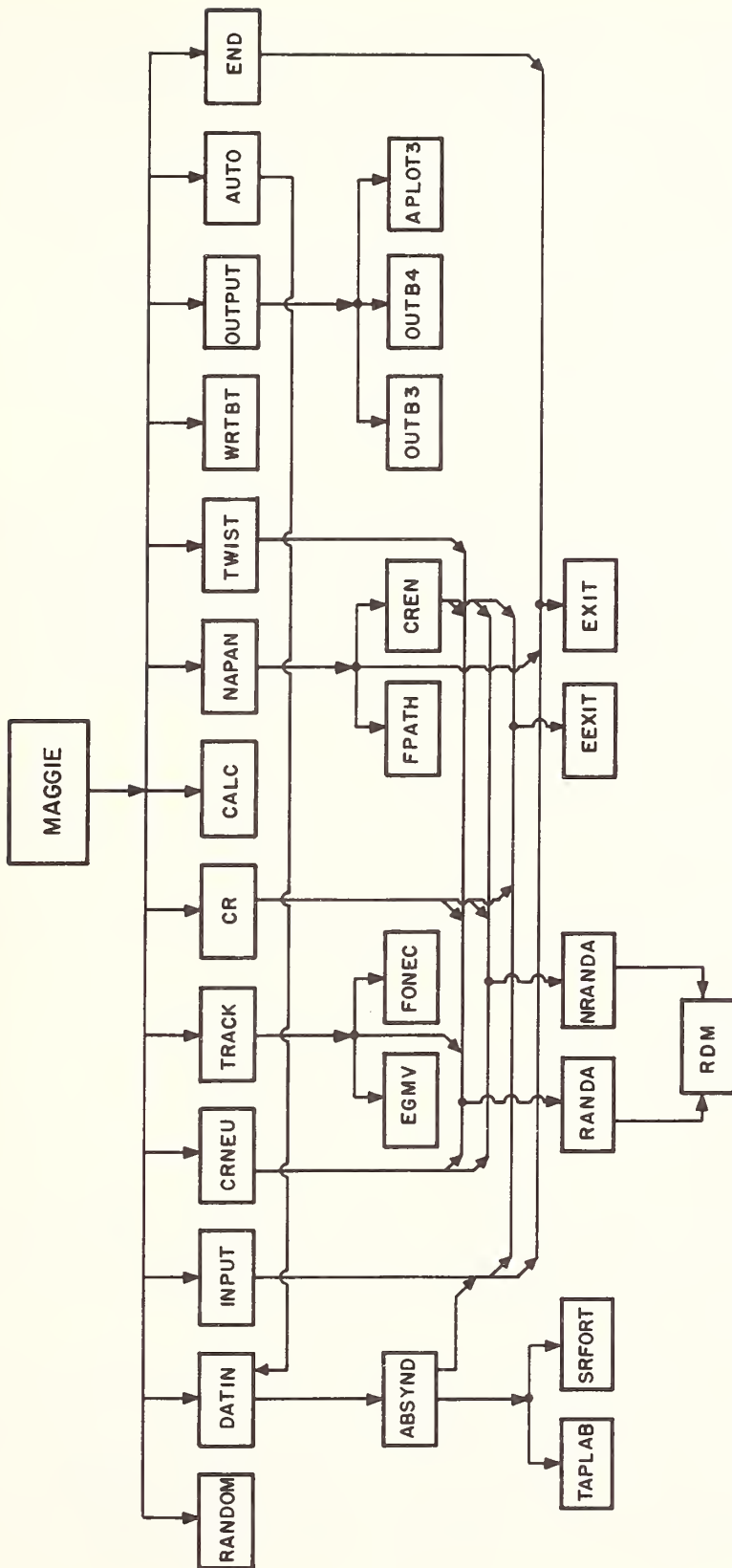


Fig. 2. The calling hierarchy for the present version of code MAGGIE.

CR, as well as c) the transmission probabilities in the direction of the detectors used in subroutines FPATH and NAPAN, are all drawn from the MOULD nuclear data file tape and not from input experimental data. The only experimental data employed are angular distributions used for the calculation of the scores for the relative probability of scattering to the various assumed detector positions in subroutine NAPAN. The elastic angular distribution (for the incident neutron energy only) used here is also automatically updated after each iteration by the subroutine AUTO. Thus it may be necessary to update the angular distributions on the MOULD file tape before and during the course of data-processing.

We now describe the changes due to the new scattering geometry in some detail for each of the subroutines affected.

1. INPUT

a. Cards 0207-0226 are changed to eliminate parameters required for samples consisting of concentric cylindrical shells, with common axes perpendicular to the beam axis, and to substitute those needed to characterize a truncated cone sample, axis lying along the neutron beam, and origin behind the neutron source. The new variables are HITE, FRAD, and ANGLE; the height, entrance face radius, and half-angle of the sample, respectively. HITE is immediately redefined as HITE/2, a more convenient quantity in subsequent calculations.

b. Card 0245, the calculation of the maximum angle subtended by the scattering sample at the neutron source, is changed to conform to the new geometry.

c. Cards 0301-0303 have been replaced. The correct expression for the flux attenuation factor for the new geometry is

$$\frac{F_o}{F} = \frac{K(o)}{\bar{K}} \cdot \frac{C}{N\lambda} \cdot \frac{h(r_1^2 + r_1\Delta r + \Delta r^2/3)}{(1 - \cos\theta_m) d^2},$$

where h is the half-length of the scattering sample, r_1 the radius of its entrance face, Δr the difference in radii of the entrance and exit faces, and d the distance from the source to the entrance face of the sample. The remaining symbols are defined as in the original.

d. The original version of MAGGIE utilized a sample which was not symmetric about the axis of the incident neutron flux. Thus the experimental sampling of scattered neutrons in the detector plane was not truly representative of the scattering into 4π , which is employed in the program to calculate the flux attenuation factors and which are used in turn to infer the integrated cross-sections. This necessitated a small correction which was calculated with the help of a classification of the outgoing

Monte Carlo tracks vs. energy and the angle with respect to the scattering sample axis. Since the present scattering sample is symmetric about the incident neutron flux the experimental sampling of scattered neutrons in the detector plane is representative, and no correction is required. Accordingly, cards 0331-0338 in INPUT, cards 1544-1548 in DATIN, and cards 1735-1743 in MAGGIE are deleted. Card 1734 of MAGGIE is replaced by the statement

20 CONTINUE

The output of this table, by cards 3695-3715 of subroutine OUTB3, is also deleted.

2. CRNEU

Cards 3117-3119, 3140-3141, 3143-3147 and 3151 are replaced as shown in the listing. The new coding creates neutrons randomly scattered over the face of the conical scattering sample, in accordance with the specified spatial distribution. The old reference to multiple materials is not deleted, but the sample is now designated "material one." Similarly, the register containing neutrons which miss the sample is retained, although for this geometry misses occur with very low frequency.

3. TRACK

The entire subroutine has been replaced. The geometry is illustrated in figure 3. The subroutine is entered with the starting point (x_o, y_o, z_o) and the direction cosines $(\cos\theta_x, \cos\theta_y, \cos\theta_z)$ already defined. A random sampling of the neutron mean free path establishes the track length to the next possible collision. The track vector is then extended until it intersects the plane determined by the endface of the sample in the direction of travel of the neutron, and the path length from the starting point to this intersection is computed. A comparison of this intersection point with the radius of the endface establishes whether the track is in the direction of the endface or the curved surface of the sample, and the program branches accordingly. If the track passes through the endface, the path length calculated above is compared with the Monte Carlo track length to the next collision and the coordinates of collision or escape from the surface of the sample as well as the time elapsed along the track are computed in a straight-forward way.

We consider now the procedure used for a neutron headed toward a curved surface of the scattering sample. The coordinates of the endpoint of the previously determined Monte Carlo track length are first found. These are used to calculate the projected distance, perpendicular to the symmetry axis, of this endpoint from the symmetry axis, as well as the length of the radius of the sample lying along this projection. Comparison of these quantities determines whether a collision occurs within the sample or not. If so, the coordinates are already known and the time elapsed along the track is all that remains to be calculated. If not, it is necessary to calculate the coordinates of the neutron's exit from

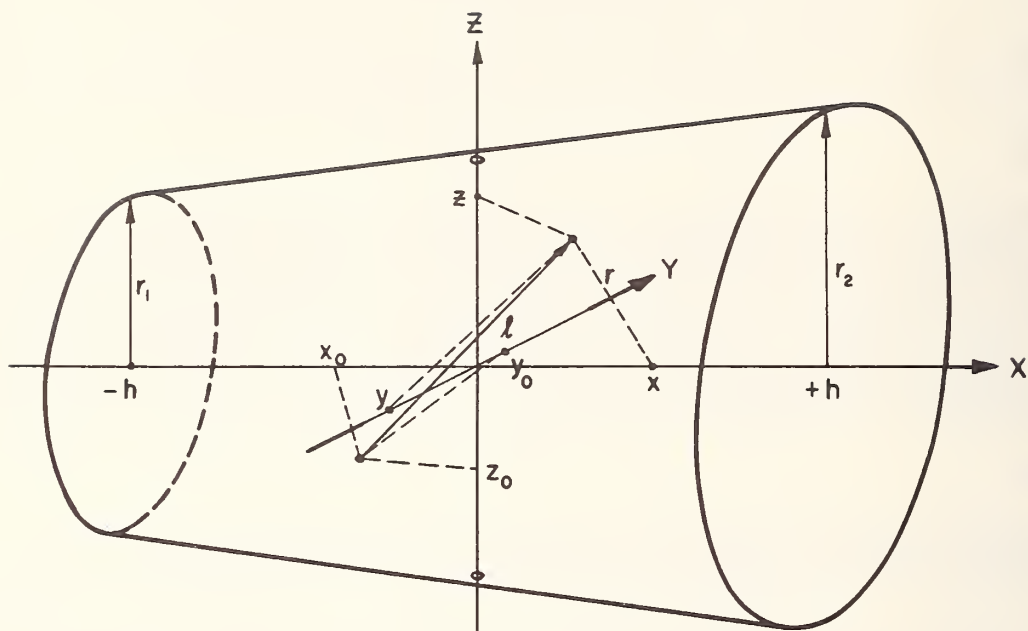


Fig. 3. The geometry for subroutine TRACK. The trajectory shown is that of a neutron which, after collision at point x_0, y_0, z_0 escapes from the curved surface of the sample at point x, y, z .

the surface of the cone. Designate ℓ the path length to the surface and x, y, z the coordinates of the intersection of the track with the surface. We have, referring to figure 3,

$$x - x_o = \ell \cos \theta_x$$

$$y - y_o = \ell \cos \theta_y$$

$$z - z_o = \ell \cos \theta_z ,$$

so that the projected distance, r , from the symmetry axis of the cone to the point (x, y, z) is given by

$$\begin{aligned} r^2 &= y^2 + z^2 \\ &= \ell^2 \sin^2 \theta_x + 2\ell(y_o \cos \theta_y + z_o \cos \theta_z) + y_o^2 + z_o^2 , \end{aligned} \quad (1)$$

where we have used the identity

$$\cos^2 \theta_x + \cos^2 \theta_y + \cos^2 \theta_z = 1 .$$

Since we know the parameters of the cone we can also calculate r from the x coordinate (again see figure 3.):

$$r = r_1 + (h+x_o) \tan \theta_o , \quad (2)$$

where θ_o is the half-angle of the cone. Eliminating r from eq'ns. (1) and (2) and arranging the terms as a quadratic in ℓ , we obtain

$$\begin{aligned} &\ell^2 \left[1 - \cos^2 \theta_x (1 + \tan^2 \theta_o) \right] \\ &+ 2\ell \left[y_o \cos \theta_y + z_o \cos \theta_z - \cos \theta_x (r_1 \tan \theta_o + h + x_o) \tan^2 \theta_o \right] \\ &+ \left[y_o^2 + z_o^2 - (r_1 + (h+x_o) \tan \theta_o)^2 \right] = 0. \end{aligned} \quad (3)$$

The positive root of this equation is the required path length to the surface, and the subroutine FONEC is then called to provide the x, y , and

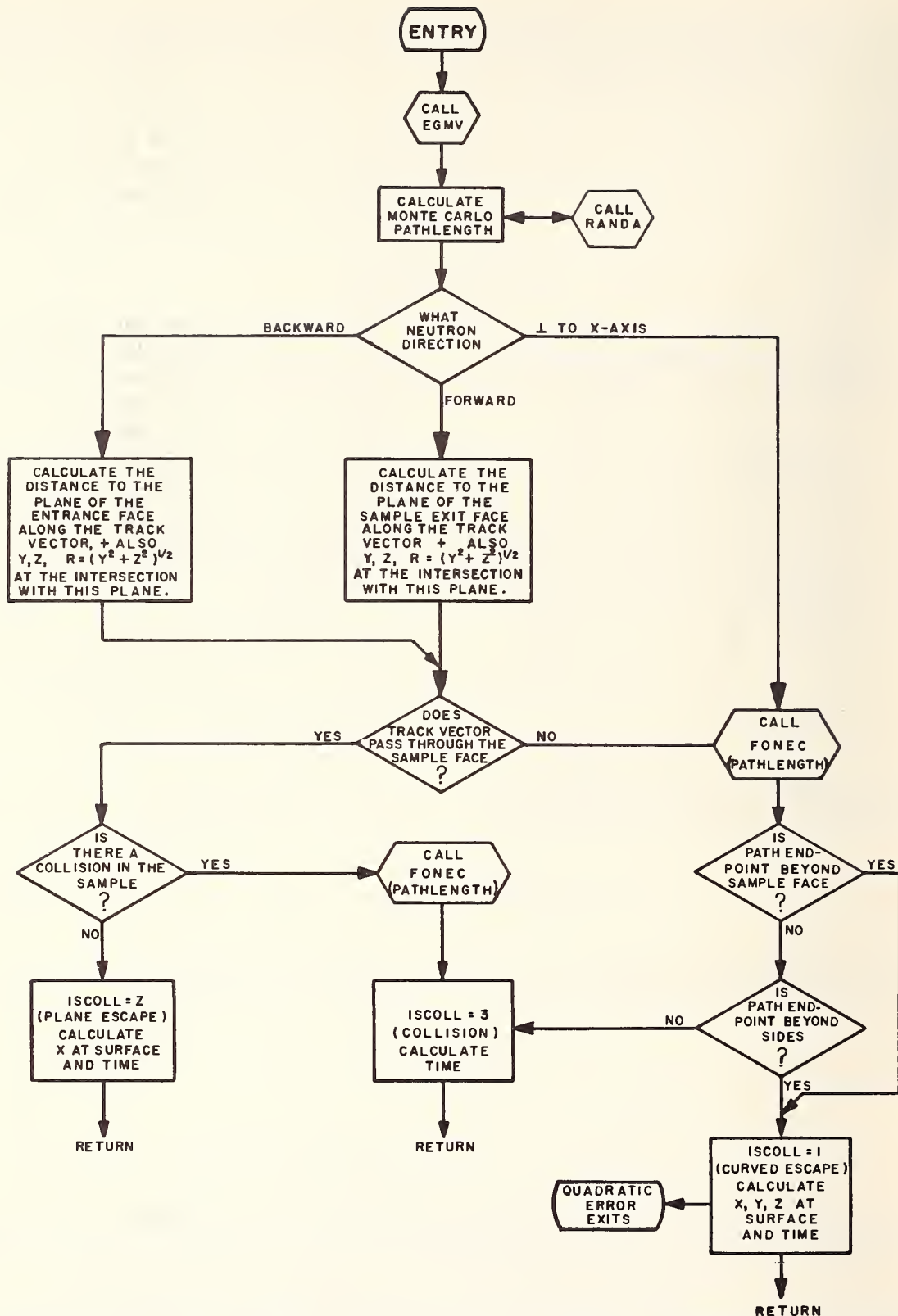


Fig. 4. Subroutine TRACK. The subroutines EGMV and FONEC return the neutron mean free path and the coordinates at the end of a track, respectively. RANDA returns a random number.

z coordinates. Error prints are provided for the cases of two positive, two negative, imaginary, or indeterminate roots as solutions to eq. (3). However none of these situations has occurred in many tens of thousands of tracks, although at least some of them are mathematically possible. Once ℓ has been calculated, determination of the coordinates of escape and the track time are straightforward. The subroutine also returns an index that indicates the fate of the neutron, i.e., curved escape, plane escape, or collision. A block diagram is shown on figure 4.

4. FPATH

This subroutine has also been entirely replaced. The same simplifying assumption is made as in the original coding--that is, the dimensions of the scattering sample are considered to be negligible compared to the flight path to the detector. This has the effect that the angle and flight path to any particular detector position may be considered to be the same from any scattering point in the sample. The quantities to be calculated are the neutron path length, p , in the sample from any point (x, y, z) to the surface of the sample and the scattering angle, θ_d , in the direction of a particular detector, at the angle ψ . The geometry is shown in figures 5 and 6. All paths are assumed to be parallel to the x-y plane, and all detector positions satisfy the condition $y \geq 0$.

We note that the locus of the intersection of a cone with a plane parallel to its axis is an hyperbola, and write, using the notation shown in figure 5,

$$\frac{X^2}{a^2} - \frac{Y^2}{b^2} = 1 \quad ,$$

$$\text{or} \quad Y^2 = X^2 \tan^2 \theta - z^2 \quad , \quad (4)$$

$$\text{since} \quad b = z \quad \text{and} \quad a = z/\tan \theta.$$

Now, referring to figure 6, we may also write

$$Y = p \sin \psi + y$$

$$\text{and} \quad X = p \cos \psi + x' \quad , \quad (5)$$

$$\text{where} \quad x' = R + h + x = \frac{r_1}{\tan \theta} + h + x.$$

Again we reduce the problem to a quadratic in path length, p , after

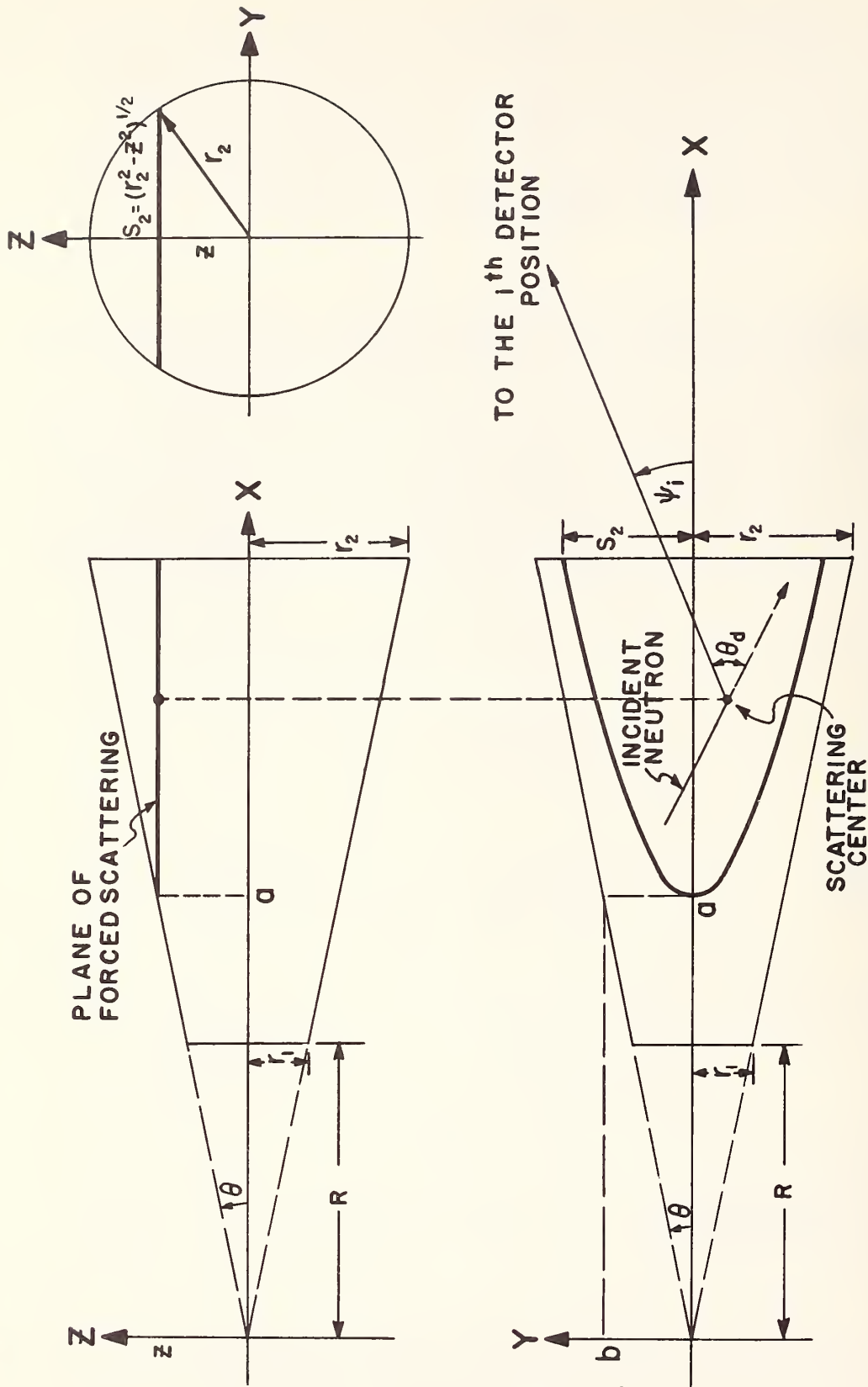


Fig. 5. The geometry for subroutine FPAT. The plane of forced scattering is located at a height $r_1 < z < r_2$ in this instance. The track of the incident neutron does not, in general, lie in this plane as shown.

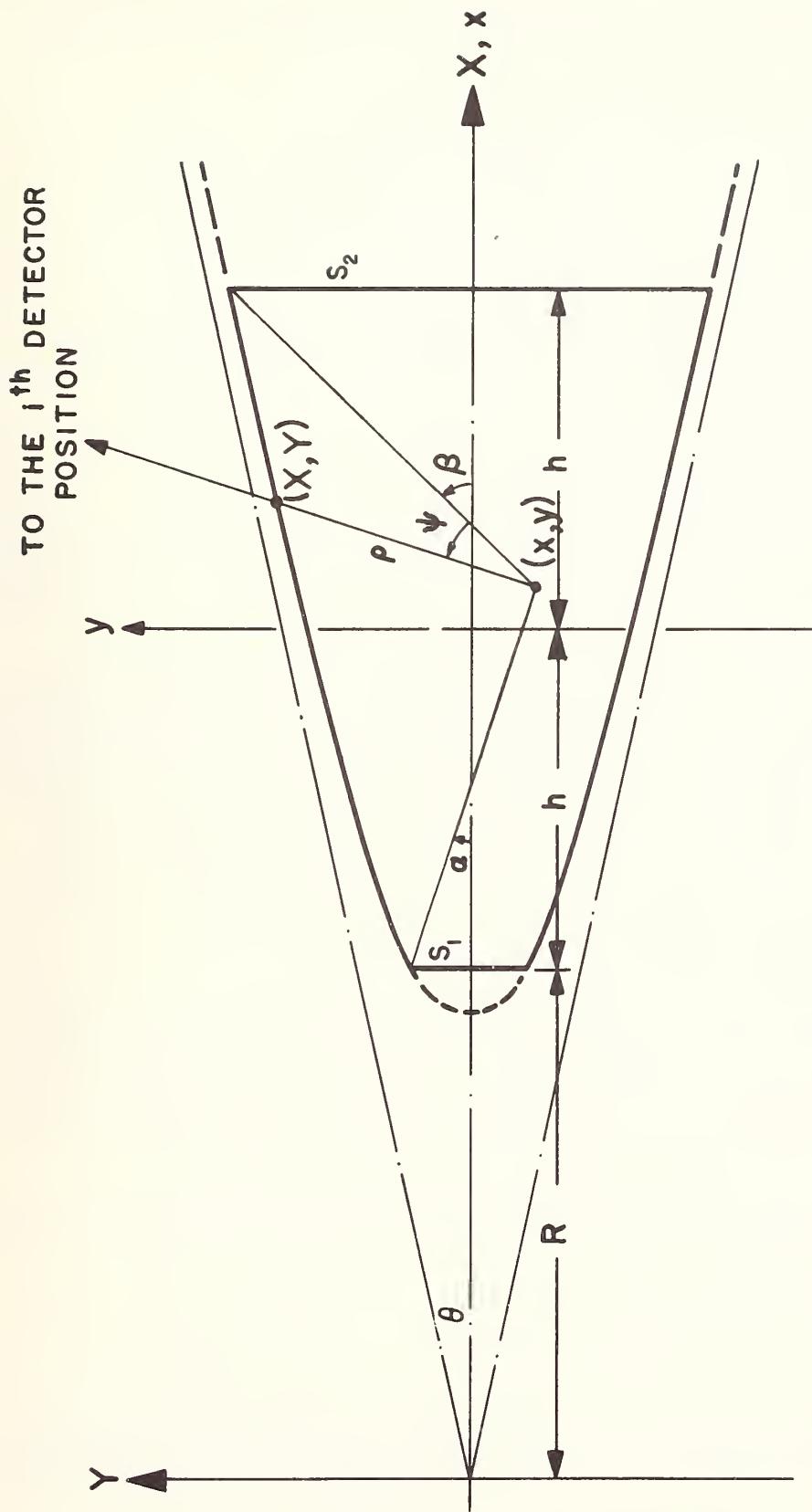


Fig. 6. A section through the scattering sample at height $z < r_1$. The upper case quantities refer to the axes X and Y , whose origin is that of hyperbola defined by the intersection of the plane (x, y) with the surface of the scattering sample.

combining eq's (4) and (5) to eliminate the variables X and Y to obtain:

$$\begin{aligned}
 & p^2(\sin^2\psi - \cos^2\psi \tan^2\theta) \\
 & + 2p(y\sin\psi - x'\cos\psi \tan^2\theta) \\
 & + (y^2 - x'^2 \tan^2\theta + z^2) = 0 .
 \end{aligned} \tag{6}$$

For each value of z there will be a pair of angles, α and β (see figure 6, which define the limits of the hyperbolic curve. These angles are given by

$$\begin{aligned}
 \tan \alpha &= \frac{(r_1^2 - z^2)^{1/2} - y}{h + x} \\
 \tan \beta &= \frac{(r_2^2 - z^2)^{1/2} - y}{h - x} .
 \end{aligned} \tag{7}$$

The subroutine is diagrammed in figure 7. After testing for the special case of exit perpendicular to the beam axis, the program tests, using eq's (7), for exit through the endfaces vs. the curved sides for the exit angle, ψ , and branches accordingly. The path length in the former case is a straightforward calculation, and for curved escape is the positive root of eq. (6). It should be noted that in the limiting case of cylinder ($\tan\theta=0$) this subroutine, unlike all of the others, does not work. A modified subroutine for a cylindrical sample is given in the Appendix following the listing of FPATH for a truncated cone.

B. Correction of Coding Errors [7]

1. INPUT

a. Statements resulting in LGR=0 in the table look-up for the mean free path (following card 0308) have all been changed to give LGR=1. LGR=0 references not a mean free path, but instead the last tabular value of collision probability in the previous MOULD table. The revised look-up supplies the mean-free path at the lowest tabulated energy for any neutron at or below that energy.

b. On the card following statement 5001, LRG has been replaced by the correct variable, LGR.

c. Card 0441 is replaced by

$$IMINM = IMAXM + 1 .$$

This change yields sequential storage of supplementary ranges of experimental angular distributions, and thus avoids loss of needed storage space in the array SUPVAL (I,J).

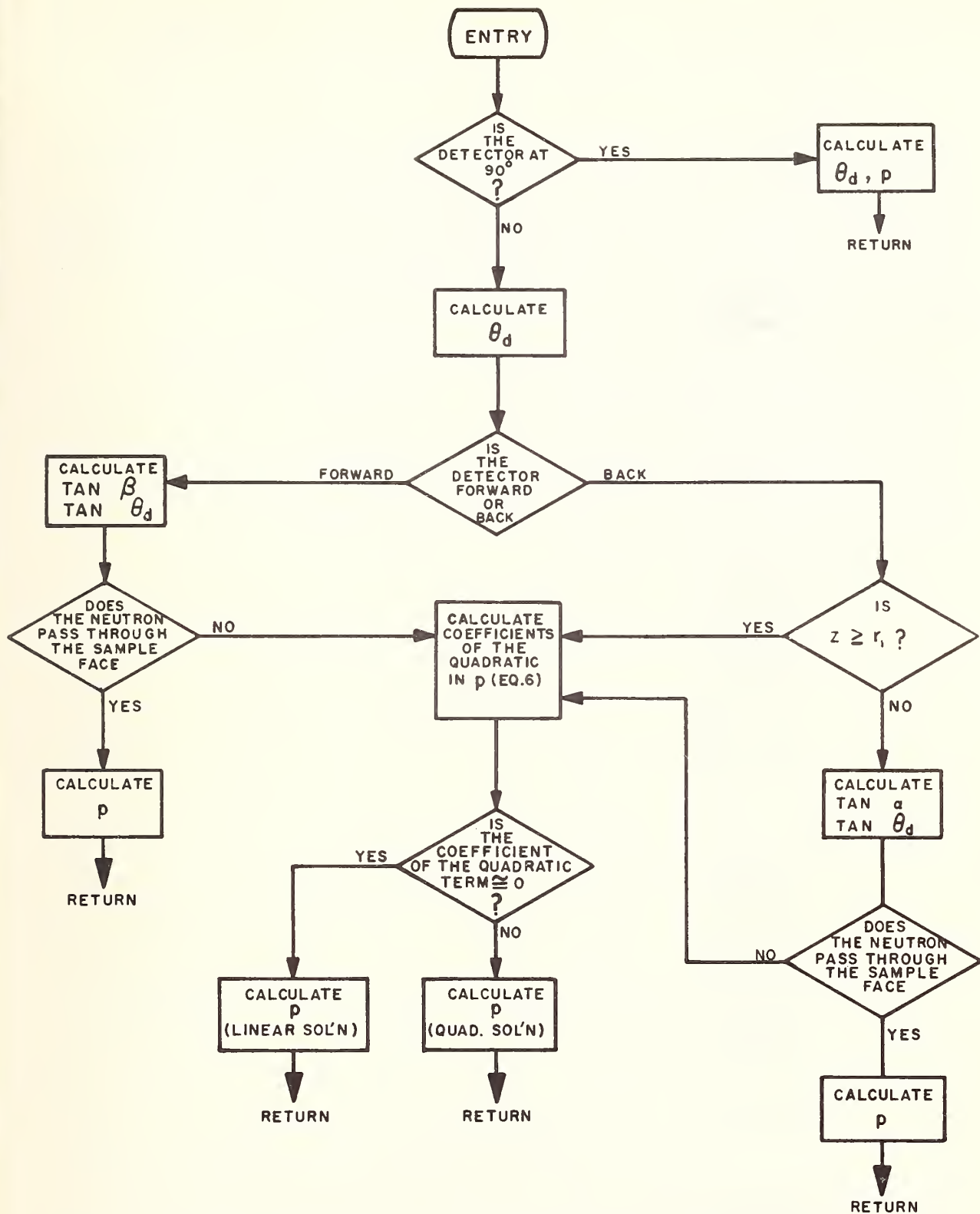


Fig. 7. Subroutine FPATH, for a truncated cone. The subroutine for a cylinder is similar.

2. EGMV

LGROUP=0 has been eliminated for the reason given in 1a. above.

3. NAPAN

a. The angle used for the interpolation of the experimental distributions, as well as for calculation of the neutron energy after collision is the laboratory angle, calculated in the subroutine FPATH. For this reason the experimental angular distributions must be read in at laboratory angles. Also the center-of-mass energy calculation on cards 3494 and 3497-3501 (in the case of center-of-mass MOULD angular distribution tables) is in error. This has been changed to a calculation in laboratory coordinates by replacing card 3494 by

```
GO TO (36,11,11,11,11,11,11,11,36,11,11,11,11,11,11) NFORM
```

and deleting cards 3497-3501.

b. The sequence of cards 3502-3512, which picks the appropriate experimental angular distribution values, does not refer to the correct angular distributions, nor to the correct neutron energy--that before collision. It has been replaced by the sequence:

```
11      IF(IMAX(J).LE.0) GO TO 15
        IMINM=IMIN(J)
        IMAXM=IMINM+IMAX(J)-1
        DO 50 JJ=IMINM,IMAXM
        IF(EIN-ENVAL(J,JJ))50,51,52

50      CONTINUE

51      KK=JJ
        GO TO 19

52      IF(JJ.LE.IMINM) GO TO 15
        KK=JJ-1
```

c. On the card following statement 5001, LRG has been replaced by the correct variable, LGR.

d. LGR=0 has been eliminated for the reason given in 1a. above.

e. The option that provides for suppression of printout of various scores, through setting the indicators IANALA, IANALB, IANALC, and IANALD negative, should not suppress storage of these scores by NAPAN. The time saved is negligible and the elastic scores, DEB(I,L), are required by the automatic iteration scheme, AUTO. Cards 3528-43 have been revised to eliminate the dependence on these printout indicators.

4. ABSYND

LGROUP=0 has been eliminated for the reason given in 1a. above.

5. CREN

For the same reason set out in section 3a. above, the center-of-mass energy calculation (cards 2960-66) has been replaced by a calculation in laboratory coordinates by deleting cards 2961-66 and replacing card 2960 with

880 GO TO 800 .

6. OUTPUT, OUTB3, OUTB4

The sum over energy, at each detector angle, stored in the vector DES(J,33) is required by subroutine AUTO. So as to allow calculation of this quantity in subroutine OUTB4, even when no B4 printout is requested, a print indicator, IFPR, has been added to the call for subroutine OUTB4. This is accomplished by the changes shown in the listing for cards 3721, 3808A, 3873, and 3975-4002.

7. AUTO

a. The quantity "chi-squared," on card 2084, should be

$$\chi^2 = \sum_1^{33} \frac{\left(\frac{\sigma(\theta)_{\text{calc}}}{\sigma(\theta)_{\text{exp}}} \exp^{-J(\theta)} \right)^2}{\sigma(\theta)_{\text{exp}}},$$

and not as formerly written, with the denominator squared.

b. It should be noted that the normalization of the Monte Carlo output on cards 2065-2068 and 2088 is not, in general, correct; but only holds for detectors positioned at equal increments of $\cos\theta$ between -1 and +1. For this reason, if subroutine AUTO is used, the input detector positions must satisfy this criterion.

8. FPATH

The original version of this subroutine contained two errors. These do not apply to the new version of FPATH reported here, but are listed for completeness. Card 3356 of the original should read

SINOM = SQRT(1-COSOM**2) ,

card 3358 be deleted, and a new card inserted following card 3359:

FP1 = WORK16**2+WORK17**2-FP2**2 .

C. Other Modifications

1. General

a. Common storage assignments have been handled by combining all common statements into one package, processed by the UNIVAC 1108 "Procedure Definition Processor" using the assembler directive FCOPY (Fortran copy). This package, designated CINCl, is then included at the time of assembly in all subroutines making use of common variables by means of the statement

INCLUDE CINCl

placed immediately following the subroutine name definition. This process, which effects a considerable economy in the size of the source deck and listing, is characteristic of the UNIVAC 1108 Assembler and FORTRAN V; however, equivalent procedures are often available to other systems. In some subroutines dimension statements still appear for those few variables not in common storage. A block of common storage has been specifically assigned to the subroutine ABSYND variables DATA and IDATA. This carries the dummy label BLANK. The details of these changes in memory storage allotment will be obvious upon examination of the listings.

b. The size of many modern computers obviates the need for linkage. We have deleted references to CHAIN(I,J) and incorporated the balance of subroutine PRELUDE into the main program MAGGIE. Similarly, references to CHAIN(I,J) in subroutines DATIN and AUTO have been replaced by the appropriate CALL and/or RETURN statements.

c. The final F in the names of all library functions (such as SIN, COS, EXP, MIN) has been deleted, as it is not compatible with FORTRAN V. In addition, such functions as FLOAT and INT have been eliminated by using the mixed expressions allowed by FORTRAN V.

d. Disc storage read and write statements have been modified in the main program MAGGIE and subroutine WRTBT to conform to 1108 FORTRAN V usage.

e. A number of formats have been altered. Some of the changes are merely different spacing options, but others are required by the new scattering geometry. A list of card numbers of the affected statements follows:

0226	0226A	0226B
0249	0258	0345
0410	0450	0451
1719	1963	1972
2111	2115	3638
3639	3828	3833

On card 2735 in subroutine CR the variable N has been replaced by the correct variable NSECS.

f. A new random number generator, subroutine RANDOM, entry RDM, has been incorporated, and the required changes in the functions RANDA and NRANDA are shown in the listings.

2. MAGGIE

a. All of the subroutine PRELUDE, except for calls to subroutines CLOCK and CHAIN, has been inserted following card 1654. A call to subroutine RANDOM immediately following the read-in of OCT, the starting value of the random number generator, initializes this generator. OCT is printed out following the call to subroutine INPUT.

b. All references to subroutine CLOCK have been deleted.

c. Card 0530 is eliminated by changing cards 1776, 1778, 1787, 1789, and 1792 to read

GO TO 1 .

3. ABSYND

a. Card 0525, which assigns the logical tape unit carrying the MOULD nuclear data tape, now reads

NUCDAT=9 .

On cards 0525, 0530, and 1394 the variable name TAPE has been changed to ATAPE to avoid confusing the compiler. These changes may not be required by other installations.

b. Cards 0553 and 0554 set all storage for the variables DATA and IDATA to zero before each ABSYND run.

c. Card 1398 is modified so as to print only those action numbers processed.

4. CR

The variable names NACT and Q have been changed to NACTV and QV, so as to avoid conflict with the array names NACT(I) and Q(I). Similarly, the variable name COS is changed to COZ, so as to make the library function COS available.

5. CREN

a. The variable names NACT and Q have been changed to NACTV and QV, for the reason given above.

b. Cards 2832-2837, 2878-2881, and 2911-2915 have been deleted, since they are not needed. In order to accomodate these deletions card 2910

has been amended to read

IF(NVCOS.LT.1) GO TO 208 ,

and card 2916 to read

$$P = (AS + \sqrt{AS^2 + A \cdot QV \cdot (1+A) / EIN + A^2 - 1}) / (1+A) .$$

6. AUTO

a. The multiple elastic sums are not required, and so cards 2069-2072 are deleted.

b. The quantity RATIO is not used, and has been eliminated from cards 2077 and 2078.

c. The do-loop 2083-2085 is redundant, since ELM2(J) is also set at card 2097. Cards 2083-2085 are therefore deleted, and card 2080 is modified to reflect this change.

d. The three do-loops in the sequence 2087-2098 have been combined into a single loop.

e. The variable name EXP(J) has been changed to EZP(J) to avoid conflict with the library function EXP.

f. Often it is desirable to separate an iteration procedure into two or more consecutive runs. Cards 2073-2082 have been modified to allow a run to be made using the partially corrected output of a previous run as input, instead of the experimental data. When this option is to be used the second field of the input card specifying the number of iterations, NTERM, should be non-zero. This is then followed by the cards specifying the experimental values for the angular distribution, FCVAL, which are removed from their usual position and replaced by the partially corrected output of the previous run.

g. In the case of rapidly varying angular distributions such as at 14.1 MeV, the usual iteration procedure is not as rapidly convergent as the "physical" method. For this method the calculated multiple scattering is first subtracted from the experimental input, and then the balance of the iteration (sample attenuation) is performed as usual by reflecting the output for single scattered neutrons about the input. The new coding on cards 2091-2097 reflects this change. For the first two iterations the iteration improves only the multiple scattering, the correction being applied both times to the experimental input. After two iterations multiple scattering is well enough known to allow the iterative procedure full play, so the correction is applied each time to the previous input, instead of the experimental values. This procedure yields much more rapid convergence, three iterations providing better convergence than six of the previous method.

The code has been applied to a typical associated-particle scattering geometry for 14.1 MeV neutrons on carbon. The scattering sample was a truncated cone of half-angle seven degrees, length 3.193 cm., and entrance face radius 2.059 cm., with its center located 20 cm from the neutron source. The experimental input used was simulated using published results for carbon at this energy. The results of a typical run are shown in Table 1, and on figures 8 and 10. The cross-section data used for the Monte Carlo scattering were those of Slaggie and Reynolds [8]. In Table 1 the quantity $\sigma(\theta)$ represents the true angular distribution, and the total and multiple outputs are the results calculated using it and the specified sample shape. The code varies $\sigma(\theta)$ until the total elastic output matches the experimental input to the accuracy required. The calculated angular distribution for neutrons undergoing multiple scattering is shown in figure 8. It can be seen that the result converges quite rapidly. Further iterations resulted only in statistical variations about the values obtained after three iterations; in fact, the largest change after two iterations, that at $\cos \theta = 0.625$, represents only a 1.5% change in the cross section.

In figure 9 we show the part of the correction, exclusive of multiple scattering, that depends upon sample shape. Also displayed for comparison is a calculation of the effect expected due to attenuation in the sample for an isotropic angular distribution. As can be seen, this correction is not independent of the input angular distribution. The quantity plotted, $\Delta\sigma/\sigma$, is the difference between the singly scattered output and the input, divided by the input, for each angle. This geometrical correction is contributed to about equally by varying path length in the sample and the change in the total neutron cross section with energy as a function of the neutron scattering angle. The correction varies from a few percent to about ten percent, and is directly dependent upon the accuracy with which the total neutron cross section is known over the range of energy exhibited by the neutron recoil. On the other hand, the correction for multiple scattering is larger, ranging from about 50% at the backward minimum to 4% at zero degrees, but depends mainly upon the experimental input data. Over most of the practical angular range of measurement it is a quite appreciable 15-20%. Figure 10 shows the input distribution used and the corrected output distribution obtained after three iterations of approximately 2000 interacting neutrons each. The total correction is, of course, largest in regions where the cross section exhibits minima. It varies from a few percent to a maximum of only twenty percent at measurable angles, since the geometrical and multiple scattering corrections tend to be in opposite directions, except at the extreme backward angles. Several examples of spectra showing elastic-inelastic and multiple inelastic effects in the energy spectra are given in the original references [4, 5]; the present version of MAGGIE also yields similar results.

TABLE 1

The scores for a typical Monte Carlo run for carbon at 14.1 MeV.

COUNTER NUMBER	Cos θ	$\sigma(\theta)$	TOTAL ELASTIC OUTPUT	MULTIPLE ELASTIC OUTPUT	EXPERIMENTAL INPUT
1	-1.0000	.023	.044	.021	.044
2	-.9375	.075	.092	.022	.093
3	-.8750	.093	.111	.024	.111
4	-.8125	.100	.120	.026	.121
5	-.7500	.110	.131	.028	.131
6	-.6875	.127	.148	.030	.148
7	-.6250	.152	.170	.031	.171
8	-.5625	.176	.190	.033	.191
9	-.5000	.202	.212	.034	.212
10	-.4375	.225	.231	.035	.232
11	-.3750	.240	.246	.036	.246
12	-.3125	.248	.253	.037	.254
13	-.2500	.245	.251	.037	.251
14	-.1875	.233	.240	.038	.241
15	-.1250	.216	.225	.038	.226
16	-.0625	.192	.205	.038	.205
17	.0000	.173	.183	.038	.183
18	.0625	.145	.155	.037	.155
19	.1250	.125	.135	.037	.135
20	.1875	.114	.125	.037	.125
21	.2500	.107	.123	.039	.122
22	.3125	.110	.127	.041	.126
23	.3750	.124	.142	.044	.140
24	.4375	.154	.170	.049	.167
25	.5000	.203	.215	.055	.212
26	.5625	.281	.285	.063	.281
27	.6250	.412	.404	.074	.400
28	.6875	.628	.605	.087	.600
29	.7500	.969	.929	.104	.925
30	.8125	1.486	1.435	.125	1.435
31	.8750	2.274	2.218	.151	2.219
32	.9375	3.451	3.380	.181	3.387
33	1.0000	5.197	5.045	.216	5.062

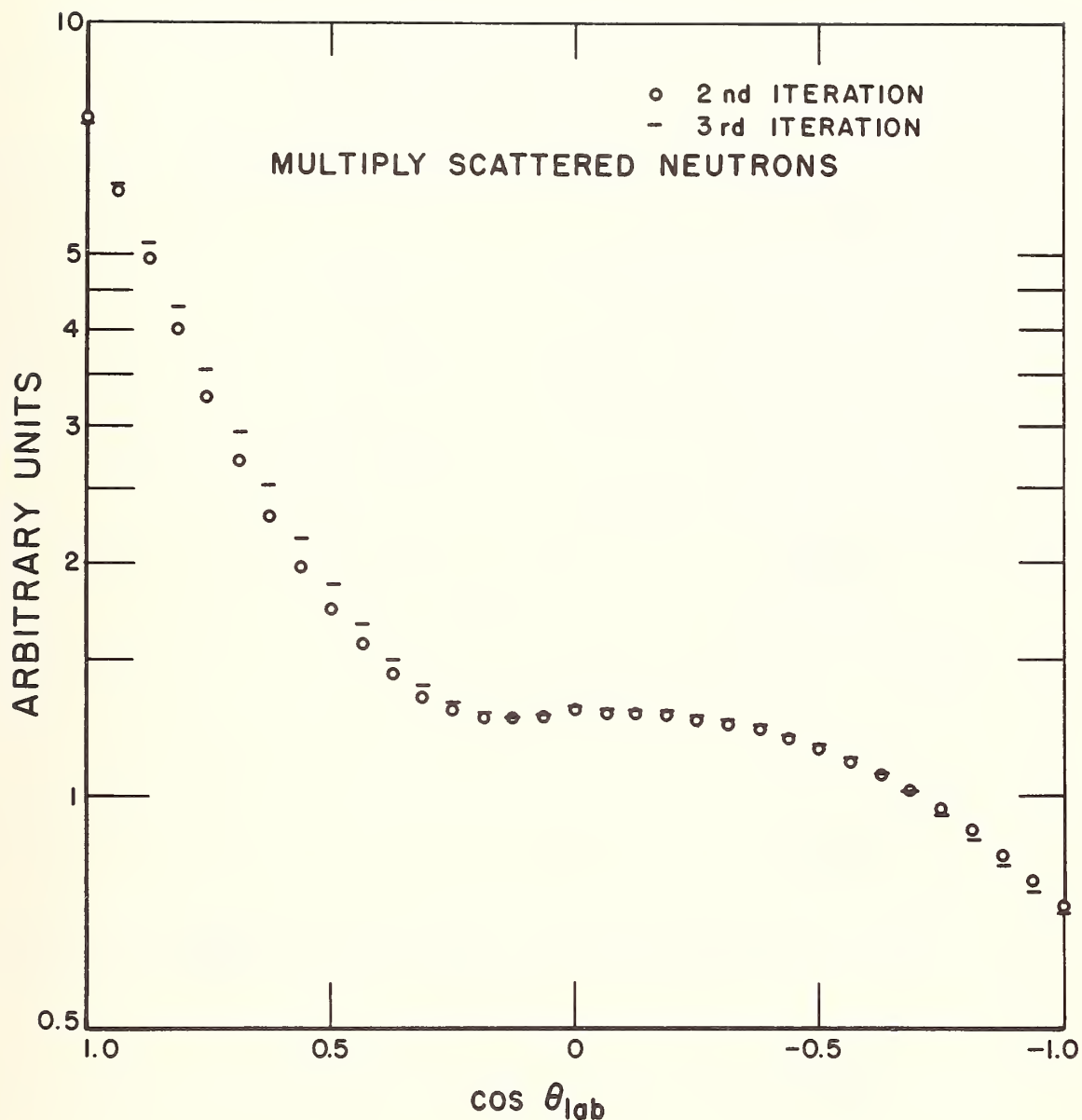


Fig. 8. The angular distribution of multiply scattered neutrons on carbon at 14.1 MeV for two successive iterations. The data for the third iteration are also given in Table 1, after suitable normalization.

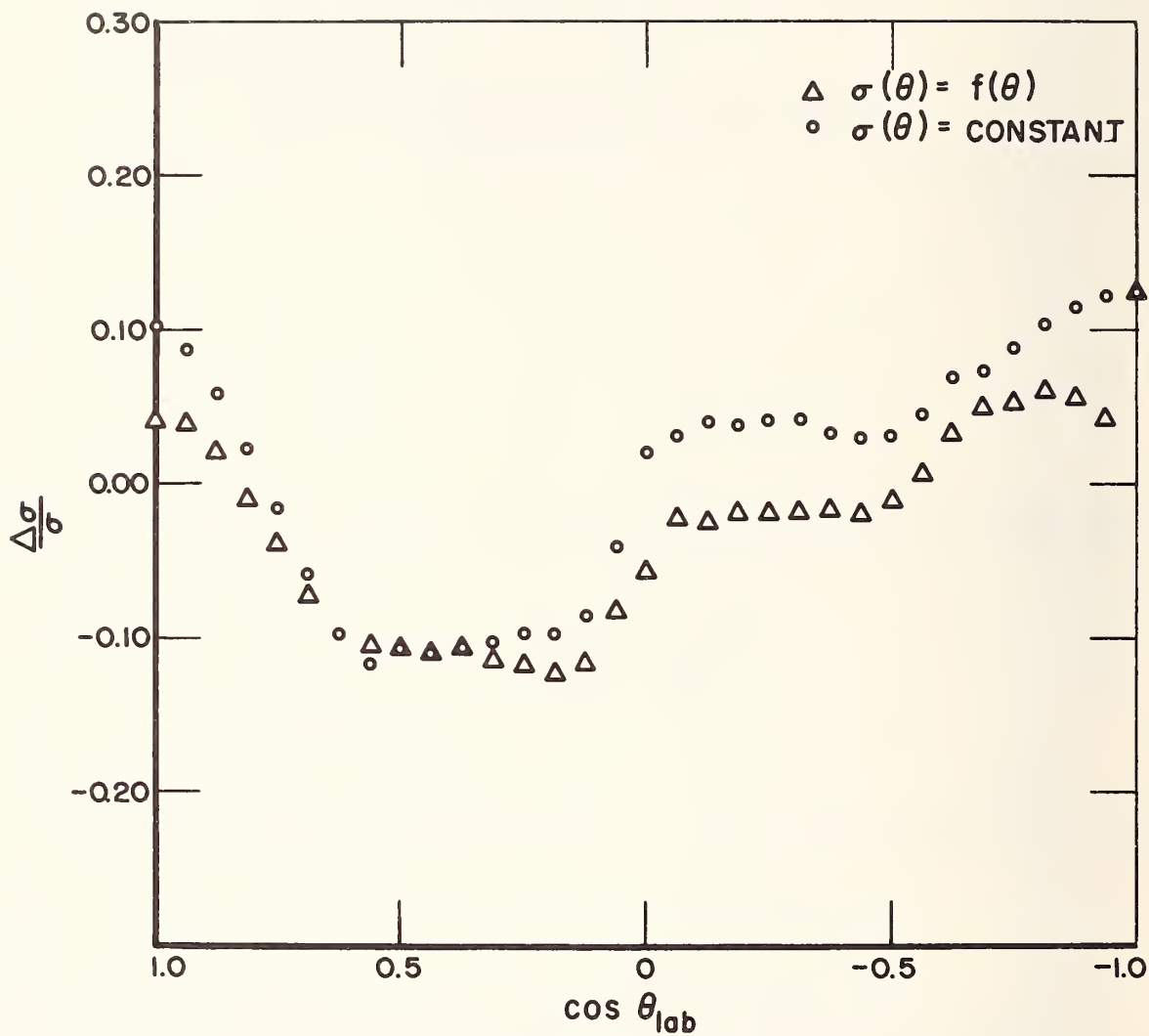


Fig. 9. The correction due to sample shape. The quantity $\Delta\sigma/\sigma$ is the change in the normalized angular distribution after subtracting multiple scattering.

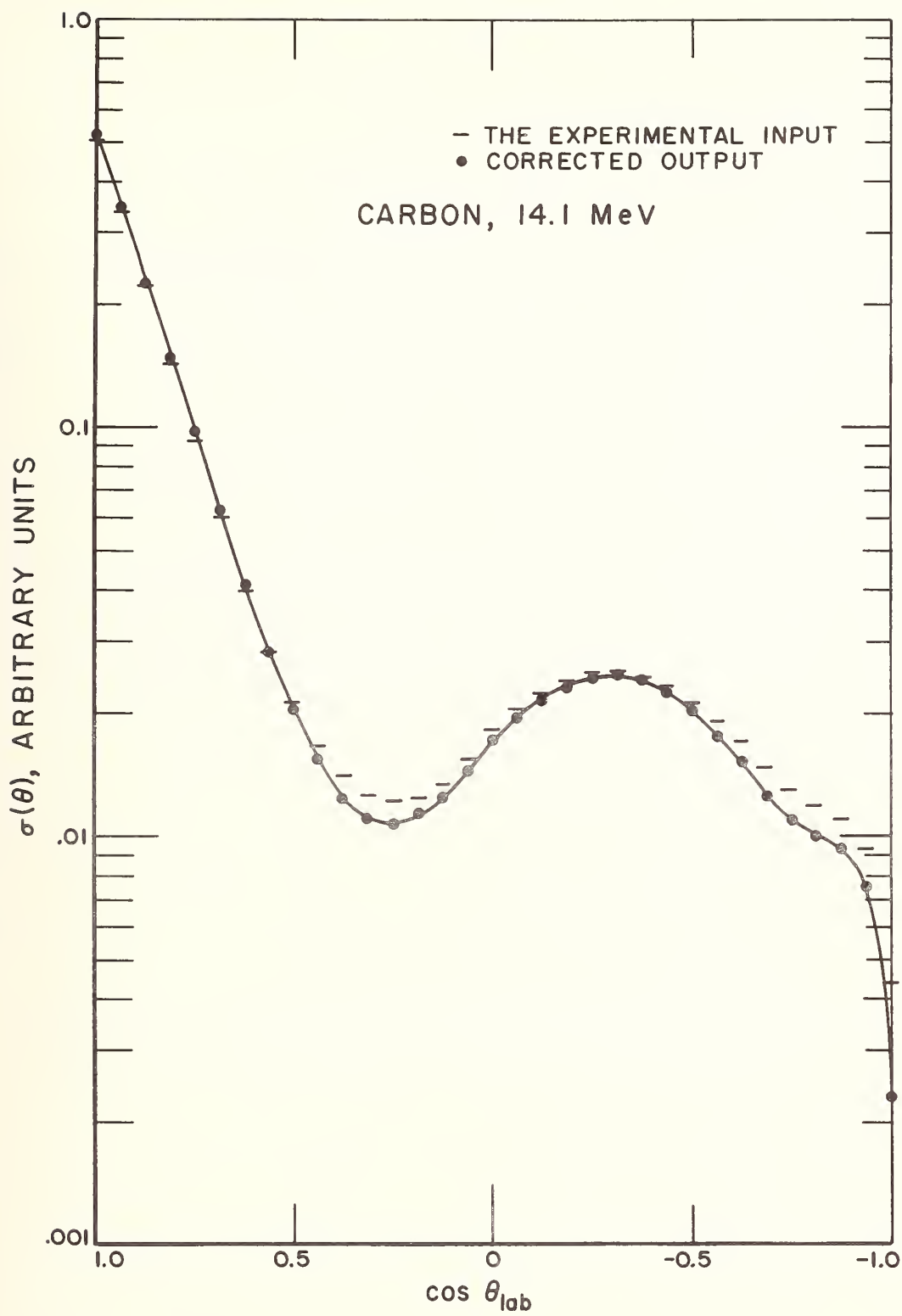


Fig. 10. The experimental input and corrected output for carbon at 14 MeV.

REFERENCES

1. Monier, L. F. C., Tripard, G. E., and White, B. L., Nucl. Instr. and Meth. 45, 282 (1966).
2. Marshak, H., Richardson, A. C. B., and Tamura, T., Phys. Rev. 150, 996 (1966).
3. A useful summary is given by M. Walt in Fast Neutron Physics, Vol. II, Marion, J. B., and Fowler, J. L., Ed. (1960).
4. Parker, J. B., Towle, J. H., Sams, D., and Jones, P. G., Nucl. Instr. and Meth. 14, 1 (1961).
5. Parker, J. B., et al, Nucl. Instr. and Meth. 30, 77 (1964).
6. Parker, K., AWRE report O-70/63; Kerr, W. M. M., AWRE report O-81/64 (1964); Miller, S. M., and Parker, K., AWRE report O-55/65.
7. Some of these coding errors have also been previously noted in private communications. We are indebted to J. B. Parker, Aldermaston, for noting the substitution of REDEN for EIN in subroutine NAPAN, as well as the errors in subroutine FPATH. H. Horstmann and H. Schmid, Geel, have noted, in addition, the difficulty involving the center of mass calculation of neutron recoil energy in subroutines NAPAN and CREN.
8. Slaggie, E. L., and Reynolds, J. T., KAPL-3099, (1966).

APPENDIX A

Listings of new and extensively revised coding.

<u>Name</u>	<u>Page</u>
CINC1	28
MAGGIE 3A	30
INPUT	36
CRNEU	44
TRACK	46
NAPAN	49
FPATH (truncated cone)	53
FPATH (cylinder)	55
OUTPUT	56
AUTO	58
RANDOM	61
RANDA	61
NRANDA	61

A complete listing is available by request from the Center for Radiation Research, Neutron Physics Section, National Bureau of Standards, Washington, D. C. 20234.

C	MAGGIE 3A	1566
C	20.7.64 A.D.PURNELL	1570
C	9.68 MODIFIED NBS VERSION FOR FORTRAN V, A.C.B.RICHARDSON	NBS1570A
C	FOR TRUNCATED CONE SCATTERING SAMPLE, INCLUDING AUTOMATIC ELASTIC	NBS1571
C	SCATTERING CALCULATION AND GRAPHS.	NBS1571A
	INCLUDE CINC1	NBS
	DIMENSION LAWTP(100)	1653
	ITER=1	1654
	INIT = 16719	NBS
	READ(5,51)UCT	NBS
	CALL RANDOM(UCT)	NBS
51	FORMAT(012)	NBS
	READ(5,200)(HEAD(I),I=1,12)	NBS
	WRITE(6,200)(HEAD(I),I=1,12)	NBS
200	FORMAT(12A6)	NBS
	CALL DATIN	NBS
	CALL INPUT	1656
	WRITE(6,52)UCT	NBS
52	FORMAT(55HUTHE STARTING VALUE FOR THE RANDOM NUMBER GENERATOR IS ,	NBS
	1012)	NBS
103	NOSAMP=0	1658
	IRECRG=0	1659
	IRGRA=0	1660
	IRGRB=0	1661
	ANEUNO=0.0	1662
	THERM=0.25000001E-7	1663
1	IF(IRGRA)2,7,4	1664
2	PRINT 3,ANEUNO	1665
3	FORMAT(1H1,47HERROR CONDITION-BIRTH REGISTER NEGATIVE ANEUNO=,F6.0	1666
	2)	1667
	IRGRA=0	1668
	IRGRB=0	1669
	GO TO 7	1670
4	IRGRA=IRGRA-11	1671
	IRGRB=IRGRB-3	1672

IWORK1=IBIRTH(IRGB+1)	1673
MAT=IBIRTH(IRGB+2)	1674
IWORK2=IBIRTH(IRGB+3)	1675
EIN=IBIRTH(IRGA+1)	1676
WORK3=BIRTH(IRGA+2)	1677
WORK4=BIRTH(IRGA+3)	1678
WORK5=BIRTH(IRGA+4)	1679
WORK6=BIRTH(IRGA+5)	1680
WORK7=BIRTH(IRGA+6)	1681
WORK8=BIRTH(IRGA+7)	1682
WORK19=BIRTH(IRGA+8)	1683
UIN=BIRTH(IRGA+9)	1684
WIN=BIRTH(IRGA+10)	1685
WORK11=BIRTH(IRGA+11)	1686
IF(IRGA-550)13,5,13	1687
IF(IIRCG)13,13,6	1688
IIRCG=IIRCG-11	1689
LOC=LOC-5600	NBS1690
READ(37)IIRCG,(BIRTH(I),I=551,4950),(IBIRTH(I),I=151,1350)	NBS1692
IRGA=IRGA+4400	1693
IRGB=IRGB+1200	1694
GO TO 13	1695
IF(ANEUNO-SAMPLE)10,8,8	1696
NOSAMP=NOSAMP+1	1697
ANEUNO=0.0	1698
CALL OUTPUT	1699
IF(ITMAG-1)81,82,81	1700
81 IF(NOSAMP-JOEFIN)10,9,9	1701
82 IF(NAC1(1)-2) 9,84,9	1702
84 IF(NOCOU-33)9,85,9	1703
85 CALL AUTO	1704
IF(ITER-10) 103,9,9	1705
9 CONTINUE	NBS1706
CALL END	1715
CALL CRNEU	1716
10 IF(NSENSE)11,13,11	1717

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11 PRINT 12 1718
12 FORMAT(120H0SERIAL STEP EN TIME FATE ISO. ACT. L M N NBS1719
13 X Y Z 1720
14 IF(NSENSE)14,15,14 1721
15 PRINT 1722
16 33,ANEUNO,IWORK1,EIN,WIN,WORK3,WORK4,WORK5 1723
17 CALL TRACK 1724
18 WORK1=WIN 1725
19 ISCOLM=ISCOLL 1726
20 GO TO (17,18,22),ISCOLM 1727
21 PESCE=PESCE+WORK1 1728
22 GO TO 18 1729
23 CESC=CESC+WORK1 1730
24 IF(NSENSE)19,20,19 1731
25 PRINT 1732
26 34,ANEUNO,IWORK1,EIN,WIN,WORK3,WORK4,WORK 1733
27 25,WORK18,WORK17,WORK18,WORK19,ISCOLM
28 20 CONTINUE
29 505 IF(IWORK1)506,506,507 1744
30 506 ANOCOL=ANOCOL+WORK1 1745
31 GO TO 1 NBS1746
32 507 IF(IWORK1-1)1,508,509 NBS1747
33 508 I=ICOLL(1) 1748
34 ONCOL(I)=ONCOL(I)+WORK1 1749
35 GO TO 1 NBS1750
36 509 NOIN=0 1751
37 DO 511 I=1,1WORK1 1752
38 IF(ICOLL(I)-2)511,511,510 1753
39 NOIN=NOIN+1 1754
40 LAOTP(NOIN)=LAOTP(I) 1755
41 CONTINUE 1756
42 IF(NOIN-2)518,512,518 1757
43 IF(LAOTP(1)-LAOTP(2))513,514,514 1758
44 I=LAOTP(2) 1759
45 J=LAOTP(1) 1760
46 GO TO 515 1761
47 514 I=LAOTP(1) 1762

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23	CALL CALC	1800
	CALL NAPAN	1801
24	WORK1=WOUT	1802
25	IF(IWORK1-100)25,25,26	1803
531	IF(NREACT(1)-2)532,531,532	1804
	ICOLL(IWORK1)=1	1805
5310	IF(IWORK1-1)26,5310,26	1806
	ALAM(1)=ALAM(1)+WORK1	1807
	GO TO 26	1808
532	IF(NREACT(1)-16)534,533,533	1809
533	ICOLL(IWORK1)=2	1810
5330	IF(IWORK1-1)26,5330,26	1811
	ALAM(2)=ALAM(2)+WORK1	1812
	GO TO 26	1813
534	DO 535 I=1,NLWREV	1814
	IF(NOLAW(1)-LAWREV(I))535,536,535	1815
535	CONTINUE	1816
	ICOLL(IWORK1)=3	1817
	LAWTP(IWORK1)=1	1818
	IF(IWORK1-1)26,5350,26	1819
5350	ALAM(3)=ALAM(3)+WORK1	1820
	GO TO 26	1821
536	ICOLL(IWORK1)=LWNO(I)+3	1822
	LAWTP(IWORK1)=LWNO(I)+1	1823
	IF(IWORK1-1)26,5360,26	1824
5360	ALAM(I+3)=ALAM(I+3)+WORK1	1825
26	IF(NSECS)32,507,27	1826
27	WORK10=WOUT	1827
	DO 31 I=1,NSECS	1828
	WORK12=LOUT(I)	1829
	WORK9=WOUT(I)	1830
	IF(WORK12-THERM)28,29,29	1831
28	WORK12=THERM	1832
	WORK9=LOG (THERM/0.25E-7)	NBS1833
29	ANGLE=COSPHI(I)	1834
	CALL TWIST(ANGLE)	1835

30	CALL WRIBT				1836
	IF(NSENSE) 30, 31, 30				1837
	PRINT				1838
	2WORK14, WORK15, WORK16, WORK17, WORK18, WORK19, ISCOLM, NONUC(I), NOACT(I)	34, ANEUNO, IWORK1, EOUT(I), UOUT(I), WOUT, WORK13,			1839
	3, NORNGE(I), NOLAW(I), NSECS				1840
31	CONTINUE				1841
32	GO TO 1				1842
33	FORMAT(1H0, F7.0, I4, 9F7.3, F10.4, I5, 3I6, I5, I7)				1843
34	FORMAT(1X, F7.0, I4, 9F7.3, F10.4, I5, 3I6, I5, I7)				1844
	END				1845

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0251 READ(5,201)NSTDP
0252 READ(5,202)(ANGS(I),STDIST(I),I=1,NSTDP)
0253 WRITE(6,2010)(ANGS(I),STDIST(I),I=1,NSTDP)
0254 FORMAT(31H0START DISTRIBUTION PROVIDED IS /9(6F10.4/))
0255 FORMAT(6I10)
0256 FORMAT(6F10.4)
0257 C DEFINE START DISTRIBUTION TO INCLUDE THE MINIMUM ANGLE WHICH WILL
0258 C JUST ENCLOSE THE SAMPLE (USUALLY ABOUT 10-16 DEGREES) BY UP TO 25 NBS0258
0259 C POINTS.
0260 DO 211 I=1,NSTDP
0261 ANGS(I)=COS (ANGS(I)/57.295828)
0262 CONTINUE
0263 DO 212 I=1,NSTDP
0264 IF(THETAM-ANGS(I))212,213,213
0265 CONTINUE
0266 WRITE(6,2120)
0267 FORMAT(41HINSUFFICIENT START DISTRIBUTION PROVIDED.)
0268 CALL EEXIT
0269 AREA(1)=0.0
0270 STDIST(I)=STDIST(I-1)-(STDIST(I-1)-STDIST(I))*(ANGS(I-1)-THETAM)/(
0271 2ANGS(I-1)-ANGS(I))
0272 ANGS(I)=THETAM
0273 L=I-1
0274 LI=I-1
0275 DO 214 K=1,L
0276 AREA(K+1)=AREA(K)+((STDIST(K+1)+STDIST(K))/2.0)*(ANGS(K)-ANGS(K+1)
0277 2)
0278 CONTINUE
0279 DO 215 K=1,64
0280 AI=K
0281 SEG=(AREA(I)*(2.0*AI-1.0))/128.0
0282 DO 216 L=1,I
0283 IF(AREA(L)-SEG)216,217,218
0284 CONTINUE
0285 WRITE(6,2160)
0286 FORMAT(63H1ERROR IN INPUT. PARTIAL SUM LRSS THAN THE TOTAL. CALL

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0287      2EXIT. )
0288      CALL EEXIT
0289      CNPHI(K)=ANGS(L)
0290      GO TO 215
0291      218  GRAD=(STDIST(L)-STDIST(L-1))/(ANGS(L)-ANGS(L-1))
0292      IF (GRAD) 300,301,300
0293      300  CNPHI(K)=ANGS(L-1)-(STDIST(L-1)-SORT ((STDIST(L-1)*STDIST(L-1))-2.NBS0293
0294      20*GRAD*(SEG-AREA(L-1)))/GRAD
0295      GO TO 215
0296      301  CNPHI(K)=ANGS(L)+(AREA(L)-SEG)/STDIST(L)
0297      215  CONTINUE
0298      WRITE(6,219) (CNPHI(I),I=1,64)
0299      219  FORMAT(55H064 EQUI-PROBABLE COSINES OF THE START DISTRIBUTION ARE/
0300      2(8F14.6))
0301      TANGLE=TAN(ANGLE)
0302      RINC=HITE*TANGLE*2.
0303      BRAD=FRAD+RINC
0304      FLUFT=(HITE*(FRAD**2+FRAD*RINC+(RINC**2)/3))/(((DIST+HITE)**2))*
0305      1(1-THETAM))
0306      M1=IXMAT+2*NMATS
0307      FNEGS=NEGS-1.0
0308      N=M1+4*NMATS
0309      MISSM=IDATA(N)
0310      UIN=LOG (STRTE/0.25E-7)
0311      S=STRTE
0312      LGR=1
0313      IF (S-EMC(1)) 5000,5000,4999
0314      4999  IF (NEGS-128) 4001,4001,5001
0315      5001  M=129
0316      LGR=M
0317      IF (S-EMC(M)) 5002,5000,5003
0318      5002  M=M-64
0319      GO TO 4020
0320      5003  M=M+64
0321      GO TO 4020
0322      4001  M=65
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0315  AVCS=(CNPH1(32)+CNPHI(33))/2.0
0316  DO 2150 I=1,LI
0317    IF(ANGS (I+1)-AVCOS)2151,2152,2150
0318  2150 CONTINUE
0319    I=LI
0320  2151 FACT=STDIST(I)-(STDIST(I)-STDIST(I+1))*((ANGS (I)-AVCOS)/(ANGS (I)
0321    2-ANGS (I+1)))
0322    GO TO 2153
0323  2152 FACT=STDIST(I)
0324  2153 FLUXFT=FLUXFT*STDIST(I)/FACT
0325    READ(5,201)NLWREV
0326    READ(5,220)(LAWREV(I),LWNO(I),I=1,NLWREV)
0327    WRITE(6,2193)
0328  2193 FORMAT(38H0LAW REFERENCE NUMBER M.C.LAW NUMBER.)
0329    WRITE(6,2194)(LAWREV(I),LWNO(I),I=1,NLWREV)
0330  2194 FORMAT(I10,I0X,I10)
0331    IF(IFB4)2192,234,2192
0332  2192 READ(5,201)IANALA,IANALB,IANALC,IANALD
0333    READ(5,201)NANAL
0334    READ(5,220) (NUCL(I),NAC1(I),I=1,NANAL)
0335  220 FORMAT(2I10)
0336    WRITE(6,221)NANAL
0337  221 FORMAT(15H0B4 ANALYSIS ON,I3,10H ACTIONS. )
0338    WRITE(6,222) (NUCL(I),I=1,NANAL)
0339  222 FORMAT(15H0NUCLIDE NOS. ,24I4)
0340    WRITE(6,223) (NAC1(I),I=1,NANAL)
0341  223 FORMAT(16H0K.P.ACTION NOS.,I3,23I4)
0342    M3=M1-NMATS
0343    IXNUC1=IDATA(M3)
0344    M4=M3+4*NMATS
0345    NUCS=IDATA(M4)
0346    DO 3 I=1,NANAL
0347    M5=IXNUC1+NUCL(I)
0348    M6=8*NUCS+M5
0349    NOACA(I)=IDATA(M6)
0350    IXACT1(I)=IDATA(M5)
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MMM=2*NOACA(I)+IXACT1(I)      0359
K=1                             0360
NOACA1=NOACA(I)                0361
DO 4 J=1,NOACA1                0362
L=MMM+K                        0363
IF(IDATA(L)-NAC1(I))7,5,7      0364
K=K+1                           0365
CONTINUE                       0366
WRITE(6,6)                     0367
CALL EXIT                      0368
NACT(I)=K                      0369
CONTINUE                       0370
FORMAT(18H K.P.ACTION NUMBER,I4,29H DOES NOT APPEAR IN THE DATA.) 0371
WRITE(6,224)                   0372
FORMAT(15H M.C.ACTION NO.,24I4) 0373
READ(5,202)(FCCOS(I),I=1,33)   0374
33 COSINES WITH 33 VALUES OF THE ANGULAR DISTRIBUTION MUST BE 0375
SUPPLIED, WITH MONOTONICALLY INCREASING VALUES OF THE COSINE.  NBS0376
DO 225 J=1,NANAL              0377
READ(5,202)(FCVAL(J,I),I=1,33) 0378
AREB=0.0                       0379
DO 228 I=1,32                 0380
AREB=AREB+ABS (FCCOS(I+1)-FCCOS(I))*((FCVAL(J,I)+FCVAL(J,I+1))/2.0NBS0381
2)                             0382
CONTINUE                       0383
IA=1                           0384
IB=15                          0385
2280 WRITE(6,226)(FCCOS(I),I=IA,IB) 0386
226 FORMAT(15H0COSINES ,15F7.3) 0387
WRITE(6,227)(FCVAL(J,I),I=IA,IB) 0388
227 FORMAT(15H DISTRIBUTION ,15F7.3) 0389
DO 229 I=IA,IB                0390
FCVAL(J,I)=FCVAL(J,I)/AREB    0391
CONTINUE                       0392
WRITE(6,230)(FCVAL(J,I),I=IA,IB) 0393
230 FORMAT(15H DIFF. X SECT. ,15F7.3) 0394

```

```

2300 IF (IA-1)2300,2300,2301
      IA=16
      IB=30
      GO TO 2280
2301 IF (IA-16)2302,2302,225
2302 IA=31
      IB=33
      GO TO 2280
225  CONTINUE
      IMINM=1
      DO 252 L=1,NANAL
      IMIN(L)=IMINM
      READ(5,250)IMAX(L)
      FORMAT(I10)
250  PRINT 251,IMAX(L),L
251  FORMAT(10H)HERE ARE,I2,46H SUPPLEMENTARY RANGES ASSOCIATED WITH ANBS0410
      ANALYSIS,I2)
253  IF (IMAX(L))252,252,253
      IMAXM=IMAX(L)-1+IMINM
      DO 254 I=IMINM,IMAXM
      READ(5,202)ENVAL(L,I)
      READ(5,202)(SUPVAL(I,K),K=1,33)
      AREB=0.0
      DO 258 J=1,32
      AREB=AREB+ABS (FCCOS(J+1)-FCCOS(J))*((SUPVAL(I,J)+SUPVAL(I,J+1))/2NBS0419
2.0)
258  CONTINUE
      IA=1
      IB=15
259  WRITE(6,2580)ENVAL(L,I)
2580 FORMAT(25H)THIS RANGE APPLIES BELOW,F7.4,5H MEV.)
      WRITE(6,226)(FCCOS(J),J=IA,IB)
      WRITE(6,227)(SUPVAL(I,J),J=IA,IB)
      DO 260 J=IA,IB
      SUPVAL(I,J)=SUPVAL(I,J)/AREB
260  CONTINUE

```

0395
0396
0397
0398
0399
0400
0401
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0430

```

261      WRITE(6,230) (SUPVAL(I,J),J=IA,IB)
      IF(IA-1)261,261,262
      IA=16
      IB=30
      GO TO 259
262      IF(IA-16)263,263,254
263      IA=31
      IB=33
      GO TO 259
254      CONTINUE
      IMIN=IMAXM+1
      IF(IMAXM-24)252,252,256
256      WRITE(6,257)
257      FORMAT(27H11MAX HAS EXCEEDED 24-HALT.)
      CALL EEXIT
252      CONTINUE
      READ(5,201)NOCOU
      READ(5,202)
      WRITE(6,232)
232      FORMAT(10H0THERE ARE,I3,56H COUNTERS LOCATED AROUND THE EQUATOR OFNBS0450
      2 THE SCATTERER AT/(10F9.4))
      NENSP
      READ(5,201)
      READ(5,202)
      WRITE(6,233)
233      FORMAT(48H0THE ENERGY SPECTRA OF THE B4 ANALYSIS CELLS ARE/(10F9.3
      2))
234      DO 2340 K=1,32
      FCCOSD(K)=FCCOS(K+1)-FCCOS(K)
2340      CONTINUE
      RETURN
      END
0431
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0440
      NBS0441
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0443
0444
0445
0446
0447
0448
0449
      NBS0450
      NBS0451
0452
0453
0454
0455
0456
0457
0458
0459
0460
0463

```

C	SUBROUTINE CRNEU	3036
C	TRUNCATED CONE SCATTERING SAMPLE, INCIDENT NEUTRONS ALONG THE AXIS	NBS
C		3116
	INCLUDE CINCL	NBS3121
61	DIMENSION TEMP(10),TRIG(4)	3122
	ITEMP=NRANDA(64)	EU.3124
62	WORK3=CNPH1(ITEMP)	EU.3126
	TEMP(1)=RANDA(-1)	3127
	TEMP(2)=RANDA(-2)	3128
	TEMP(3)=TEMP(1)*TEMP(1)	3129
	TEMP(4)=TEMP(2)*TEMP(2)	3130
	TEMP(5)=TEMP(3)+TEMP(4)	3131
	IF(TEMP(5))62,62,63	3132
63	IF(TEMP(5)-1.0)64,62,62	3133
64	TRIG(1)=(TEMP(3)-TEMP(4))/TEMP(5)	NBS3134
	TRIG(2)=(2.0*TEMP(1)*TEMP(2))/TEMP(5)	3135
	TRIG(3)=SQRT(1.0-WORK3*WORK3)	3136
	WORK4=TRIG(1)*TRIG(3)	3137
	WORK5=TRIG(2)*TRIG(3)	3138
C	DOES THE NEUTRON HIT THE TARGET	3139
C		NBS3140
C	TEMP(1)=FRAD+(DIST+HITE)*TRIG(3)/WORK3	3142
	IF(TEMP(1))67,67,65	NBS3143
65	TEMP(2)=-((DIST+HITE)/WORK3	NBS3145
66	WORK6=-HITE	NBS3146
	WORK7=WORK4*TEMP(2)	NBS3147
	WORK8=WORK5*TEMP(2)	3148
	ANEUNO=ANEUNO+1.0	3149
	EIN=STRTE	NBS3150
	UIN=LOG(EIN/0.25E-7)	NBS3151
	MAT=1	3152
	IWORK1=0	3153
	WORK19=0.0	3154
	WIN=1.0	

3155
3156
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3160

67
WORK11=0.0
IWORK2=0
RETURN
AMISS=AMISS+1.0
GO TO 61
END

SUBROUTINE TRACK

C TRACKS NEUTRONS IN A TRUNCATED CONE WITH AXIS LYING ALONG THE NEUTRON BEAM

```

INCLUDE CINC1
CALL EGMV
PATHL=(-(LOG(RANDA(1))))*PATH
IF(WORK3)1,6,2
1 TRS=-(HITE+WORK6)/WORK3
  RAD2=FRAD**2
  GO TO 3
2 TRS=(HITE-WORK6)/WORK3
  RAD2=BRAD**2
3 WORK17=WORK7+TRS*WORK4
  WORK18=WORK8+TRS*WORK5
  TRR2=WORK17**2+WORK18**2
  IF(RAD2.LT.TRR2)GO TO 6
  IF(TRS.GT.PATHL)GO TO 13
  ISCOLL=2
  IF(WORK3.GT.0)GO TO 4
  WORK16=-HITE
  GO TO 5
4 WORK16=HITE
5 WORK19=WORK19+TRS/SPEED
  RETURN
6 CALL FONEC(PATHL)
  IF(HITE.LE.ABS(WORK16))GO TO 7
  TRA2=WORK17**2+WORK18**2
  RAD2=(FRAD+(HITE+WORK16)*TANGLE)**2
  IF(RAD2.GT.TRA2)GO TO 14
7 ISCOLL=1
  NRTS=0
  RAD=FRAD+(HITE+WORK6)*TANGLE
  TA=1.-WORK3**2*(1.+TANGLE**2)
  TB=WORK7*WORK4+WORK8*WORK5-WORK3*TANGLE*RAD

```

```

TC=WORK7**2+WORK8**2-RAD**2
IF(ABS(TA).GT.1.E-30)GO TO 8
IF(ABS(TB).LT.1.E-30)GO TO 21
TR=-TC/(2.*TB)
GO TO 12
8 TD=TB*TB-TA*TC
IF(TD)15,10,9
9 TE=-TB+SQRT(TD)
IF(TE.LT.0)GO TO 10
NRTS=1
TR=TE/TA
10 TE=-TB-SQRT(TD)
IF(TE.LT.0)GO TO 11
NRTS=NRTS+1
TR=TE/TA
11 IF(NRTS-1)17,12,19
12 CALL FONEC(TR)
WORK19=WORK19+TR/SPEED
RETURN
13 CALL FONEC(PATHL)
14 ISCOLL=3
WORK19=WORK19+PATHL/SPEED
RETURN
C ERROR EXITS
15 WRITE(6,16)
16 FORMAT(24H IMAGINARY ROOT IN TRACK)
GO TO 23
17 WRITE(6,18)
18 FORMAT(26H NO POSITIVE ROOT IN TRACK)
GO TO 23
19 WRITE(6,20)
20 FORMAT(28H TWO POSITIVE ROOTS IN TRACK)
GO TO 23

```

```
21 WRITE(6,22)
22 FORMAT(28H INDETERMINATE ROOT IN TRACK)
23 WRITE(6,24) WORK3,WORK4,WORK5,WORK6,WORK7,WORK8,PATHL
24 FORMAT(7F10.3)
   CALL EXIT
   END
```

SUBROUTINE NAPAN	3371
INCLUDE CINCI	
DIMENSION FACT(33),PATHLL(33),ANGLEE(33),KL(33)	
NENSP1=NENSP-1	NBS3451
M1=IXMAT+NMATS+MAT	3453
N=M1+4*NMATS	3454
MISSM=IDATA(N)	3455
M2=IDATA(M1)-MISSM	3456
DO 1 I=1,NOCOU	3457
COSOM=COUNT(I)	3458
CALL FPATH(COSOM,PATHL,ANGLEF)	3459
DO 2 K=1,32	NBS3460
IF(FCCOS(K+1)-ANGLEF)2,3,3	3461
CONTINUE	NBS3462
K=32	3463
AFCOS=FCCOS(K)	3464
FACT(I)=(ANGLEF-AFCOS)/FCCOSD(K)	3465
PATHLL(I)=PATHL	NBS3466
ANGLEE(I)=ANGLEF	3467
KL(I)=K	NBS3468
CONTINUE	3469
DO 7 J=1,NANAL	3470
NFORM=0	3471
NAC=NAC1(J)	3472
NUCLID=NUCL(J)	3473
PART=PART1(J)	3474
ATOM=ATOM1(J)	3475
WSQ=ATOM*ATOM-1.0	3476
DO 6 I=1,NOCOU	3477
K=KL(I)	3478
AS=ANGLEE(I)	3479
FACTOR=FACT(I)	3480
PAT:IL=PATHLL(I)	3481
ASQ=AS*AS	3482
IF(PART)7,7,8	3483
	3484

```

8      IF(NAC-2)9,10,9
10     X=SQRT (ASQ+WSQ)
      P=((AS+X)/(ATOM+1.0))*((AS+X)/(ATOM+1.0))
      REDEN=EIN*P
      GO TO 11
9      IF(NFORM)34,34,35
34     CALL CREN
      REDEN=EINER
      IF(REDEN)38,11,11
35     GO TO (36,11,11,11,11,11,11,11,11,11,11,11,11,11),NFORM
36     P=(AS+SQRT (ASQ+ATOM*QV*(1.0+ATOM)/EIN+WSQ))/(1.0+ATOM)
      REDEN=EIN*P*P
11     IF(IMAX(J).LE.0)GO TO 15
      IMINM=IMIN(J)
      IMAXM=IMINM+IMAX(J)-1
      DO 50 JJ=IMINM,IMAXM
      IF(EIN-ENVAL(J,JJ))50,51,52
50     CONTINUE
51     KK=JJ
52     IF(JJ.LL.IMINM)GO TO 15
      KK=JJ-1
      GO TO 19
15     SIGMA=(FCVAL(J,K+1)-FCVAL(J,K))*FACTOR+FCVAL(J,K)
      GO TO 16
19     SIGMA=(SUPVAL(KK,K+1)-SUPVAL(KK,K))*FACTOR+SUPVAL(KK,K)
16     DO 20 L=1,NENSP1
      IF(ENSP(L+1)-REDEN)20,21,21
20     CONTINUE
      L=NENSP1
21     S=REDEN
      LGR=1
      IF(S-EMC(1))5000,5000,4999
4999   IF(NEGS-128)4001,4001,5001
5001   M=129
      LGR=M

```

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NBS3486
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NBS3494
NBS3495
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NBS3503
NBS3504
NBS3505
NBS3506
NBS3507
NBS3508
NBS3509
NBS3510
NBS3511
NBS3512
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NBS EU.
EU.
EU.
EU.NBS

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```

SUBROUTINE FPATH(COSOM,PATHL,ANGLEF)

C   PATHLENGTH AND ANGLE FOR B4 ANALYSIS, FOR TRUNCATED CONE SAMPLE.
C   AXIS ALONG THE INCIDENT NEUTRON BEAM.

    INCLUDE CINC1

C   NEUTRON OUT AT 90 DEGREES.
    IF (ABS(COSOM).GE.1.E-5)GO TO 1
    ANGLEF=WORK4
    TSQ=TANGLE**2
    XPR=FRAD/TANGLE+HITE+WORK16
    PATHL=SQRT(XPR*XPR+TSQ-WORK18**2)-WORK17
    RETURN

C   NEUTRON OUT ENTRANCE FACE
    1 SINOM=SQRT(1.-COSOM**2)
    ANGLEF=COSOM*WORK3+SINOM*WORK4
    IF (COSOM.GT.0)GO TO 2
    IF (FRAD.LT.ABS(WORK18))GO TO 3
    FTAN=(SQRT(FRAD**2-WORK18**2)-WORK17)/(HITE+WORK16)
    TANOM=SINOM/COSOM
    IF (FTAN+TANOM.LT.0)GO TO 3
    PATHL=-(HITE+WORK16)/COSOM
    RETURN

C   NEUTRON OUT EXIT FACE
    2 BTAN=(SQRT(BRAD**2-WORK18**2)-WORK17)/(HITE-WORK16)
    TANOM=SINOM/COSOM
    IF (BTAN.LT.TANOM)GO TO 3
    PATHL=(HITE+WORK16)/COSOM
    RETURN

C   NEUTRON OUT THE SIDE, BUT NOT AT 90 DEGREES, NORMAL SOLUTION.
    3 TSQ=TANGLE**2
    XPR=FRAD/TANGLE+HITE+WORK16

```

```

PA=SINOM**2-(COSOM**2)*TSQ
PB=WORK17*SINOM-XPR*COSOM*TSQ
PC=WORK17**2-(XPR**2)*TSQ+WORK18**2
IF(ABS(PA).LT.1.E-30)GO TO 4
PATHL=(SQRT(PB**2-PA*PC)-PB)/PA
RETURN

```

```

C      LINEAR SOLUTION
4      PATHL=-PC/(2.*PB)
      RETURN
      END

```

```

SUBROUTINE FPATH(COSOM,PATHL,ANGLEF)
PATHLENGTH AND ANGLE FOR THE B4 ANALYSIS, FOR A CYLINDRICAL SAMPLE
WITH ITS AXIS ALONG THE INCIDENT NEUTRON BEAM.

INCLUDE CINC1
PATHL=SQRT(FRAD**2-WORK18**2)-WORK17
IF(ABS(COSOM).GT.1E-5)GO TO 1
ANGLEF=WORK4
GO TO 3

1 SINOM=SQRT(1-COSOM**2)
TANOM=SINOM/COSOM
ANGLEF=COSOM*WORK3+SINOM*WORK4
IF(COSOM.GT.1E-5)GO TO 2
FTAN=PATHL/(HITE+WORK16)
IF((FTAN+TANOM).LT.1E-5)GO TO 3
PATHL=- (HITE+WORK16)/COSOM
RETURN

2 FTAN=PATHL/(HITE-WORK16)
IF(FTAN.LT.TANOM)GO TO 3
PATHL=(HITE+WORK16)/COSOM
RETURN

3 PATHL=PATHL/SINOM
RETURN
END

```

C	SUBROUTINE OUTPUT	3878
	OUTPUTS RESULTS AT THE END OF EACH SAMPLE.	3879
	INCLUDE CINC1	NBS
	DIMENSION P(33)	3958
	PRINT	3960
1	1,(HEAD(I),I=1,12)	3961
	FORMAT(12A6)	3962
2	PRINT	3963
	2,NOSAMP	3964
	FORMAT(1H0,52X,13HSAMPLE NUMBER,I3)	3965
3	PRINT	3966
	3,SAMPLE	3967
	FORMAT(1H0,40X,F8.0,33H NEUTRONS STARTED IN EACH SAMPLE.)	3968
	PART =NOSAMP	3969
	TOT=PART*SAMPLE	3970
	PRINT	3971
	5,TOT	3972
5	FORMAT(1H0,39X,F9.0,35H NEUTRONS HAVE BEEN TRACKED SO FAR.)	3973
	PRINT	3974
	4,PESC,CESC	NBS3975
4	FORMAT(24HUPANE SURFACE ESCAPES =,F11.3/24H0CURVED SURFACE ESCAPE	NBS3937
	2S=,F11.3)	NBS3978
	CALL OUT83	NBS3979
	I1=ENSP-1	NBS3980
	IF(IFB4.EQ.0)GO TO 7	3982
	CALL OUT84(DEA,1,IANALA)	NBS3984
	IF(IANALA.LE.0)GO TO 8	NBS3985
	DO 6 I=1,I1	NBS3986
	P(I)=DEA(IANALA,I)	NBS3987
6	CALL APL0T3(P,I1,ENSP,IANALA)	3989
8	CALL OUT84(DEB,2,IANALB)	NBS3991
	IF(IANALB.LE.0)GO TO 10	NBS3992
	DO 9 I=1,I1	NBS3993
9	P(I)=DEB(IANALB,I)	NBS3994
	CALL APL0T3 (P,I1,ENSP,IANALB)	3996
10	CALL OUT84(DEC,3,IANALC)	
	IF(IANALC.LE.0)GO TO 12	
	DO 11 I=1,I1	
11	P(I)=DEC(IANALC,I)	
	CALL APL0T3 (P,I1,ENSP,IANALC)	


```

12 CALL OUIB4(JED,4,IANALD)
   IF(IANALD.LE.0)GO TO 7
   DO 13 I=1,I1
13  P(I)=DED(IANALD,I)
   CALL APLOTS (P,I1,ENSP,IANALD)
7   RETURN
   END

```

```

NBS3998
NBS3999
NBS4000
NBS4001
   4003
   4004
   4005

```

SUBROUTINE AUTO	1977
C TO ALLOW AUTOMATIC ITERATION UNDER THE FOLLOWING CONDITIONS.	1978
C NUMBER OF COUNTERS AND INPUT POINTS = 33.	1979
C ONE NUCLIDE ONLY.	1980
C ELASTIC CORRECTION ONLY.	1981
C PHYSICAL ITERATION.	
C (DELTA SIGMA)/(SIGMA) IS THE CHANGE IN THE CROSS SECTION EXCLUDING	
C MULTIPLE SCATTERING.	
INCLUDE CINC1	NBS
DIMENSION TOTEL(33),CHI(33),ELM2(33),XIP3(33)	NBS2062
SUMEL=(DEB(1,33)+DEB(33,33))/32.0	2065
DO 4 J=2,32	2066
SUMEL=SUMEL+DEB(J,33)/16.0	2067
4 CONTINUE	2068
IF(ITER.NE.1)GO TO 500	NBS2073
READ(5,101)NTERM,IFCONT	NBS2077
101 FORMAT(2I10)	NBS2078
NTERM=NTERM-1	2079
IF(IFCONT.EQ.0)GO TO 100	NBS
READ(5,104)(EZP(I),I=1,33)	NBS
104 FORMAT(6F10.4)	NBS
SUMEZP=(EZP(1)+EZP(33))/32	NBS
DO 107 J=2,32	NBS
107 SUMEZP=SUMELZP+EZP(J)/16	NBS
DO 108 J=1,33	NBS
108 EZP(J)=EZP(J)/SUMEZP	NBS
GO TO 102	NBS
100 DO 105 I=1,33	2074
EZP(1)=FCVAL(1,I)	NBS2075
105 CONTINUE	2076
102 IF(NTERM.LE.1)GO TO 300	NBS2080
106 READ(5,103)(SAMP(I),I=1,NTERM)	NBS2081
103 FORMAT(6F10.2)	2082
300 CHISO=0.0	2086

```

DO 303 J=1,33
TOTAL(J)=DEB(J,33)/SUMEL
CHI(J)=((TOTAL(J)-EZIP(J))*2)/EZIP(J)
CHISQ=CHISQ+CHI(J)
DEDNJ=DED(J,33)/SUMEL
ELM2(J)=FCVAL(1,J)
FACTOR=((TOTAL(J)-DEDNJ)/ELM2(J)
IF(ITER.GT.2)GO TO 302
XIP3(J)=(EZIP(J)-DEDNJ)/FACTOR
GO TO 303
302 XIP3(J)=FCVAL(1,J)/FACTOR
303 CONTINUE
214 SUM=(XIP3(1)+XIP3(33))/32.0
DO 207 J=2,32
SUM=SUM+XIP3(J)/16.0
207 CONTINUE
DO 208 I=1,33
FCVAL(1,I)=(XIP3(I))/SUM
208 CONTINUE
DO 220 I=1,33
220 XIP3(I)=DEB(I,33)-DED(I,33)
SUMSEL=(XIP3(1)+XIP(33))/32
DO 221 I=2,32
221 SUMSEL=SUMSEL+XIP3(I)/16
DO 222 I=1,33
222 XIP3(I)=(XIP3(I)/SUMSEL-ELM2(I))/ELM2(I)
PRINT 209,ITER
209 FORMAT(1H1,20H AFTER ITERATION NO.,I3,49H . NORMALISATION IS UNIT
1AREA FOR ELASTIC EVENTS.)
216 ITER1=ITER+1
PRINT 210,ITER,ITER,ITER,ITER1
210 FORMAT(1H0,7HCOUNT,4X,6HCOSSINE,11X,3HEXP,7X,10HOUTPUT NO.,I2,4X,
13HNO.,I2,5H CHI ,7X,3HNO.,I2,11X,3HNO.,I2,4X,11HDELTA SIGMA,2X, NBS2112
27HCOUNTER/2X,6HNUMBER,5X,5HANGLE,10X,5HINPUT,7X,11HALL ELASTIC,6X, 2113
37HSQUARED,8X,5HINPUT,11X,5HINPUT,9X,6H/SIGMA,2X,6HNUMBER) NBS2114
PRINT 111,(J,COUNT(J),EZIP(J),TOTAL(J),CHI(J),ELM2(J),FCVAL(1,J), NBS2115

```

NBS2087
2088
NBS2089
2090
NBS
2097
NBS
NBS
NBS
NBS
NBS
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2099
2100
2101
2102
2103
2104
2105
NBS
NBS
NBS
NBS
NBS
NBS
NBS
2106
2107
2108
2109
2110
2111
NBS2112
2113
NBS2114
NBS2115

```

      XIP3(J),J,J=1,33)
111  FORMAT(14,7F15.6,I4)
      PRINT 112,CHISQ
112  FORMAT(1H0,8HCHISQ = ,F10.6)
997  IF(ITER-NTERM)415,416,416
415  SAMPLE=SAMP(ITER)
      ITER=ITER+1
      GO TO 999
416  ITER=10
      GO TO 998
999  CALL DATIN
998  RETURN
      END

```

```

NBS2116
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```


A Brief Description of the Components of MAGGIE

1. MAGGIE (Main Program). This program calls the various subroutines required for the analysis, retrieves neutrons from disc storage when required, outputs track parameters if desired, and records the various fates of tracked neutrons.
2. RANDOM (and RDM). This subroutine contains the random number generator, RDM, as a separate entry. RANDOM is called at the beginning of program MAGGIE to enter the starting value for RDM.
3. DATIN serves to set all of the output arrays to zero, and calls ABSYND.
4. ABSYND reads the required nuclear data from the MOULD tape and puts it in encoded sequential storage in the array DATA-IDATA for use during the Monte Carlo tracking.
5. INPUT reads and processes samples and experimental angular distribution data from card input and calculates most of the flux attenuation factor.
6. CRNEU creates random incident neutrons at the entrance face of the sample, in accordance with the input source distribution.
7. TRACK tracks neutrons in the sample, specifying coordinates of collision or escape.
8. EGMV. This subroutine computes the mean free path, velocity, and lethargy group number.
9. FONEC calculates coordinates at the end of a track from initial position, direction cosines, and track length.
10. CR. This subroutine, using random sampling of the information stored by ABSYND from the MOULD data tape, determines all of the parameters of a collision.
11. CALC determines some constants used in subroutine NAPAN.
12. NAPAN. This subroutine calculates and scores the probability of detection at each detector for each collision in the sample.
13. FPATH calculates the path length in the direction of each detector for each collision.
14. CREN is an abbreviated version of CR used by NAPAN that determines only the neutron energy.

15. TWIST chooses new direction cosines after a collision.
16. WRIBT stores any secondary neutrons produced, for recall at the end of the current tracking.
17. OUTPUT, OUTB3, OUTB4 and APLOT3. These subroutines print the results of the calculation.
18. AUTO performs the calculations required for iterative correction of the elastic angular distribution, prints the current output, and calls DATIN to begin the next iteration.
19. END, and its entries EXIT and EEXIT designate normal vs. error exit conditions.
20. TAPLAB returns the tape logical unit label.
21. SRFORT is a subroutine for skipping tape records.

Input requirements for MAGGIE-NBS

	<u>FORMAT</u>	<u>VARIABLE</u>	<u>COMMENTS</u>
1.	012	OCT	Octal starting value for RDM.
2.	12A6	HEAD(I)	Arbitrary heading. Column one should be a 1.
3.	110	NMATS	The number of materials in the sample. This is always = 1.
4.	110, E10.4	IDATA(MAT5)	Number of different nuclides in the material.
		DATA(MAT4)	Density of the material.
5.	110, E10.4	IDATA(K2)	Nuclide reference number (i.e. position on the MOULD tape: see output of MOULD for this).
		DATA(K4)	Proportion of this nuclide in the material. This card is repeated for each nuclide in the material.
6.	2I10	IFB4	Positive for B4 output.
		NSENSE	Positive for track print.
7.	3F10.4	HITE	Length of the samples, in cm.
		FRAD	Entrance radius of same, in cm.
		ANGLE	Half-angle of same, in radians.
8.	2I10	JOBFIN	Number of independant samples (= 1 if ITMAG > 0).
		ITMAG	Positive for automatic iteration.
9.	F10.4	SAMPLE	Number of neutrons in each sample.
10.	F10.4	STRTE	Starting energy in MeV. If the starting energy is set = 0 a fission spectrum is assumed.
11.	F10.4	DIST	Distance from the center of the sample to the source (negative), in cm.
12.	110	NSTDP	No. of start distribution points to be read in (≤ 25).
13.	6F10.4	ANGS(I)	Three pairs to a card, the angle and start distribution for that angle. The start distribution need not be normalized, and this card is repeated until NSTDP pairs are read in.

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|--------------------------------------------------------------|--------|--------------------------------------|----------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------|
| 14. | I10 | NLWREV | The number of Law Reference numbers for inelastic scatter, i.e. the number of Law Reference numbers to different angular distributions on the MOULD output (P.C.N.'s* 4-15 only) (NLWREV \leq 50). |
| 15. | 2I10 | LAWREV(I)
LWNO(I) | One pair per card, the above Law Reference numbers and the monotonically increasing "Monte Carlo" law numbers allocated. Any inelastic law not given here will be printed in the B3 section results under Law zero. |
| NOTE: The following cards are not required if IFB4 \leq 0. | | | |
| 16. | 4I10 | IANALA
IANALB
IANALC
IANALD | Four markers, for the tables: A) Complete multiple scatter analysis. B) Elastic events only. C) Inelastic events only. D) Multiple elastic events only. A negative marker suppresses the table. A positive marker = N produces, in addition to the table, a graph for the Nth counter. |
| 17. | I10 | NANAL | Number of actions (index I below) (i.e. different neutron processes, such as elastic, inelastic from the 1st excited state, etc.) to be processed by the sub-routine NAPAN. |
| 18. | 6F10.4 | NUCL(I)
NAC1(I) | The nuclide reference number (as in card 5), and the P.C.N.* for each action. There are NANAL such pairs. |
| 19. | 6F10.4 | FCCOS(J) | Thirty-three values of cosine, including -1 and +1, monotonically increasing. |
| 20. | 6F10.4 | FCVAL(I,J) | Thirty-three values of the angular distribution (index J) corresponding to the above values of cosines. Repeat for each action (index I) in the same order as in 18. |

NOTE: The following three items are repeated, as a group, for each action.

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|-----|-----|---------|--------------------------------------------------------------------------|
| 21. | I10 | IMAX(I) | Number (index I) of supplementary ranges for this action (even if zero). |
|-----|-----|---------|--------------------------------------------------------------------------|

*The "P.C.N.'s (Particular Classification Numbers) are listed starting on p.8 of AWRE report no. O 70/63, "The Aldermaston Nuclear Data Library as at May, 1963", K. Parker.

NOTE: The following two items are repeated for each supplementary range. If, for any action, there is no supplementary range these items are omitted.

22.	F10.4	ENVAL(I, L)	Upper limit of the range in MeV. These must be listed in order of decreasing value.
23.	6F10.4	SUPVAL(L, J)	Thirty-three values (index J) of the distribution (range L) at the cosine values FCCOS(J).
24.	I10	NOCOU	Number of counters (≤ 33).
25.	6F10.4	COUNT(I)	Cosines of counter locations.
26.	I10	NENSP	Number of output energy points (≤ 33). The output will be classified into the bins formed by these points.
27.	6F10.6	ENSP(I)	Values of output energy points in MeV.
28.	2I10	NTERM IFCONT	Number of iterations (≤ 10). If this is a continuation of a previous run, for further iterations this should be non-zero (see Section II.c.6.f).
29.	6F10.4	SAMP(I)	Number of neutrons for each iteration, <u>except</u> the first.

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